

# Second-Order Sufficient Optimality Conditions for Optimal Control of Non-Smooth, Semilinear Parabolic Equations

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Non-smooth and Complementarity-based Distributed Parameter Systems: Simulation and Hierarchical Optimization

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# SECOND-ORDER SUFFICIENT OPTIMALITY CONDITIONS FOR OPTIMAL CONTROL OF NON-SMOOTH, SEMILINEAR PARABOLIC EQUATIONS

#### LIVIA BETZ\*

Abstract. This paper is concerned with an optimal control problem governed by a non-smooth, semilinear parabolic PDE. The nonlinearity in the state equation is only directionally differentiable, locally Lipschitz continuous, and is allowed to have infinitely many non-differentiable points. By employing its limited properties, Bouligand-differentiability of the control-to-state map is shown. This enables us to establish second-order sufficient optimality conditions. We provide concrete settings where these reduce to the first-order necessary optimality condition.

Key words. Optimal control of PDEs, non-smooth optimization, second-order sufficient conditions

## AMS subject classifications. 49J20, 35K58, 49K99

1. Introduction. In this paper we establish second-order sufficient conditions for the following optimal control problem:

$$\begin{array}{ccc}
\min_{u \in L^{r}(0,T;L^{2}(\Omega))} & J(y,u) \\
\text{s.t.} & \dot{y}(t) + A y(t) + f(y(t)) = B u(t) \quad \text{a.e. in } (0,T) \\
& y(0) = 0,
\end{array}$$
(P)

where  $\Omega \subset \mathbb{R}^n$ , n = 2, 3, is a bounded Lipschitz domain, J is a smooth function, A is a linear unbounded operator and f is a non-smooth mapping. The precise statements will be given at the end of this section and in Assumption 2.1 below. The essential feature of (P) is that the nonlinearity f appearing in the state equation is only directionally differentiable. Thus, the second-order analysis cannot be performed by classical techniques for smooth optimization problems in Banach spaces.

Optimal control problems subject to non-smooth constraints are challenging even in the finite dimensional case, see e.g. [29] and the references therein. Difficulties arise from the non-smoothness of the control-to-state mapping, which does not allow to apply the standard Karush-Kuhn-Tucker (KKT) theory. For this reason, various optimality conditions of different strength have been introduced, such as e.g. Clarke (C), Bouligand (B), and strong stationarity. In the spirit thereof, stationarity concepts for the infinite dimensional case are defined in [18]. The most rigorous stationarity concept is strong stationarity. In a previous work [23], necessary optimality conditions of this type were established, from which we will benefit in the present paper.

While second-order sufficient optimality conditions (SSC) for the optimal control of smooth PDEs have been intensively investigated, see e.g. [4, 5, 7-9, 11, 14, 26, 28] and the references therein, the literature on SSC for the optimal control of non-smooth *problems* is rather rare. To the best of our knowledge, the only contributions in this field deal with elliptic VIs. These were addressed in [21] (obstacle problem) and [3] (static elastoplasticity). In [24] it was proven that the obstacle control problem is

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convex if the desired state is behind the obstacle. This result was extended in [21], where sufficient conditions for the optimal control of the obstacle problem in the general case were presented. To the best of the author's knowledge, the investigation of second-order sufficient conditions for optimization problems governed by non-smooth parabolic PDEs is an open research topic.

What distinguishes the problem (P) from the ones analyzed in [3] and [21] is not only the parabolic component but also the very general non-smooth mapping f. For example, we allow the set of non-differentiable points of f to be at most countable. We require (in addition to strong stationarity, positive-definiteness/coercivity of the Hessian of the "Lagrangian") only a sign condition on the adjoint state and we can provide settings where this is fulfilled. By comparison, in [21] sign conditions are imposed not only on the adjoint state but also on the multiplier, while in [3] regularity assumptions on the adjoint state and multipliers are made. If the nonlinearity f is twice continuously differentiable, then the SSC derived in the present work coincide with the classical second-order sufficient optimality conditions, see Remark 4.19 below.

The paper is organized as follows. In Section 2 we state the precise assumptions on the data and lay the foundations for our analysis, by recalling some crucial results from [23]. Section 3 is mainly devoted to proving the Bouligand-differentiability of the control-to-state operator in appropriate spaces. This is the first essential step towards the investigation of SSC, which will be performed in Section 4. The latter contains our main results, namely two sets of second-order sufficient conditions for the optimal control of (P) (Theorems 4.15 and 4.18 below). In Section 5 we present a setting where the necessary optimality condition alone is sufficient for optimality.

**Notation.** Throughout the paper, C and c denote generic positive constants. If X and Y are two linear normed spaces, the space of linear and bounded operators from X to Y is denoted by  $\mathcal{L}(X,Y)$ . For the open ball in X around  $x \in X$  with radius R we write  $B_X(x,R)$ . The symbol  $X^*$  stands for the dual space of X, while  $\langle ., . \rangle_X$  stands for the dual pairing between X and  $X^*$ . If X is compactly embedded in Y, we write  $X \hookrightarrow Y$ , and  $X \stackrel{d}{\hookrightarrow} Y$  means that X is dense in Y. If X and Y are Banach spaces, we use the notation  $[X, Y]_{\theta}$  for the complex and  $(X, Y)_{\theta,\omega}$  for the real interpolation space, respectively, where  $\theta \in (0, 1)$  and  $\omega \in [1, \infty]$ , see e.g. [30]. If a linear operator A is the infinitesimal generator of a semigroup, the latter will be denoted by  $\{e^{tA}\}_{t\geq 0}$ , see also [25, Chp. 2.5]. In all what follows, T > 0 is a fixed final time and  $\Omega \subset \mathbb{R}^n$ , n = 2, 3, is a bounded Lipschitz domain in the sense of [22, Chp. 1.1.9]. For simplicity, we abbreviate  $Q := (0, T) \times \Omega$ . The boundary of  $\Omega$  consists of two disjoint measurable parts  $\Gamma_D$  and  $\Gamma_N$ . By  $W_D^{1,q}(\Omega)$  we denote the closure of the set  $\{\psi|_{\Omega} : \psi \in C_0^{\infty}(\mathbb{R}^n), \operatorname{supp}(\psi) \cap \Gamma_D = \emptyset\}$  with respect to the  $W^{1,q}$ -norm, where  $q \in (1,\infty)$ . For the dual space associated with  $W_D^{1,q'}(\Omega)$  we use the symbol  $W^{-1,q}(\Omega)$ , where q' stands for the conjugate exponent. Moreover, we abbreviate

$$\begin{split} \mathbb{W}_{0}^{r}(W_{D}^{1,q}(\Omega),W^{-1,q}(\Omega)) &:= \{ v \in W^{1,r}(0,T;W^{-1,q}(\Omega)) \cap L^{r}(0,T;W_{D}^{1,q}(\Omega)) : v(0) = 0 \}, \\ \mathbb{W}_{T}^{r'}(W_{D}^{1,q'}(\Omega),W^{-1,q'}(\Omega)) &:= \{ v \in W^{1,r'}(0,T;W^{-1,q'}(\Omega)) \cap L^{r'}(0,T;W_{D}^{1,q'}(\Omega)) : v(T) = 0 \}, \\ \text{where } r \in (1,\infty). \end{split}$$

2. Standing assumptions and known results. This section is devoted to collecting the assumptions on the data as well as crucial results from [23] concerning the state equation.

ASSUMPTION 2.1. For the quantities in (P) we require the following:

1. Let

$$n < q < \infty$$
 and  $\frac{2q}{q-n} < r < \infty$  (2.1)

be fixed.

- 2. The operator  $A: W^{-1,q}(\Omega) \to W^{-1,q}(\Omega)$  is linear, unbounded, and closed. Its domain of definition is given by  $W_D^{1,q}(\Omega)$ . In addition,  $0 \notin \sigma(A)$ , where  $\sigma(A)$  denotes the spectrum of A.
- 3. Moreover, A satisfies maximal parabolic  $L^r(0,T;W^{-1,q}(\Omega))$ -regularity, i.e., for every  $g \in L^r(0,T;W^{-1,q}(\Omega))$ , the equation  $\dot{w} + Aw = g$  admits a unique solution  $w \in W^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega))$ .
- The nonlinearity f : R → R is assumed to be monotone increasing. Moreover, it is Lipschitz continuous on bounded sets, i.e., for all M > 0, there exists a constant L<sub>M</sub> > 0 such that

$$|f(z_1) - f(z_2)| \le L_M |z_1 - z_2| \quad \forall \, z_1, z_2 \in [-M, M].$$

5. The function f is directionally differentiable at every point, i.e.,

$$\left|\frac{f(x+\tau h) - f(x)}{\tau} - f'(x;h)\right| \stackrel{\tau \searrow 0}{\longrightarrow} 0 \quad \forall x, h \in \mathbb{R}.$$

Moreover, the set of non-differentiable points of f is at most countable.

- 6. The operator  $B: L^2(\Omega) \to W^{-1,q}(\Omega)$  is linear and bounded.
- 7. The objective  $J : L^{r}(0,T; W^{1,q}_{D}(\Omega)) \times L^{r}(0,T; L^{2}(\Omega)) \to \mathbb{R}$  is twice continuously Fréchet-differentiable.

REMARK 2.2. With a little abuse of notation, the Nemytskii-operators associated with f and B, considered with different ranges, will be denoted by the same symbol.

Comments regarding the above assumptions are provided at the end of this section. From now on, Assumption 2.1 is tacitly assumed in the following without mentioning it every time. We point out that this fits in the general setting of [23, Assumptions 2.1 and 2.5], as explained in the following. With the notations from the preceding contribution [23] we have

$$X = W^{-1,q}(\Omega), \quad \mathcal{D} = W^{1,q}_D(\Omega), \quad U = L^2(\Omega), \quad Y = L^{\infty}(\Omega).$$

Note that the maximal parabolic regularity assumption on A implies that A generates an analytic semigroup, see [2, Section 3]. The Nemytskii operator  $f : L^{\infty}(\Omega) \to L^{\infty}(\Omega)$  is well-defined and Lipschitz continuous on bounded sets, i.e., for every M > 0, there exists  $L_M > 0$  so that

$$\|f(y_1) - f(y_2)\|_{L^{\infty}(\Omega)} \le L_M \|y_1 - y_2\|_{L^{\infty}(\Omega)} \quad \forall y_1, y_2 \in B_{L^{\infty}(\Omega)}(0, M).$$
(2.2)

When considered with domain  $L^{\infty}(\Omega)$  and range  $L^{\beta}(\Omega)$ ,  $\beta < \infty$ , f is directionally differentiable, see proof of [23, Lemma 6.4] for details. Therein, Assumption 2.1.1 on r and q is justified as well. The condition q > n guarantees that there exists  $\theta \in (0,1)$  such that  $(W^{-1,q}(\Omega), W_D^{1,q}(\Omega))_{\theta,\infty} \hookrightarrow L^{\infty}(\Omega)$ , see [23, (6.4)], while the relation between r and q in (2.1) ensures that  $r(1-\theta) > 1$ , i.e., [23, (2.4)]. These turn out to be essential not only for the existence of solutions of the state equation, see [23], but also for the upcoming second-order analysis, see Remark 2.6 below. Note that, in view of [30, Thm. 1.15.2 (d), p.101] and [25, Thm. 2.6.13, p.74], it holds

$$\|e^{-tA}\|_{\mathcal{L}(W^{-1,q}(\Omega),L^{\infty}(\Omega))} \le c t^{-\theta} \quad \forall t \in (0,T],$$

$$(2.3)$$

which will be crucial in the proof of Theorem 3.2 below. Let us mention that we dropped the density assumption on B and the convexity assumption on J, as they were needed in [23] just for deriving (strong stationary) necessary optimality conditions and for proving the existence of global minimizers, respectively, which is not the case in this paper.

The following embedding will be crucial in the next sections and is a consequence of [1, Thm. 3] combined with  $W_D^{1,q}(\Omega) \hookrightarrow W^{-1,q}(\Omega), (W^{-1,q}(\Omega), W_D^{1,q}(\Omega))_{\theta,1} \hookrightarrow (W^{-1,q}(\Omega), W_D^{1,q}(\Omega))_{\theta,\infty} \hookrightarrow L^{\infty}(\Omega)$ , see [1, Section 3]:

$$\mathbb{W}_{0}^{r}(W_{D}^{1,q}(\Omega), W^{-1,q}(\Omega)) \hookrightarrow \hookrightarrow C([0,T]; L^{\infty}(\Omega)).$$

$$(2.4)$$

Since the setting in Assumption 2.1 is just a special case of [23, Assumption 2.1], we can apply the general results in [23, Sections 2-3] on our state equation. We begin by introducing the control-to-state mapping.

DEFINITION 2.3. The solution operator of

$$\dot{y}(t) + Ay(t) + f(y(t)) = Bu(t) \quad a.e. \text{ in } (0,T),$$
  
$$y(0) = 0$$
(2.5)

is denoted by  $S: L^r(0,T;L^2(\Omega)) \ni u \mapsto y \in W^r_0(W^{1,q}_D(\Omega),W^{-1,q}(\Omega))$ . Note that, in view of [23, Proposition 2.11], this is well-defined.

PROPOSITION 2.4. [23, Proposition 2.11] The control-to-state mapping S is Lipschitz continuous on bounded sets, i.e., for every R > 0, there exists a constant  $L_R > 0$  such that, for all  $u_1, u_2 \in \overline{B_{L^r(0,T;L^2(\Omega))}(0,R)}$ , it holds

$$\|S(u_1) - S(u_2)\|_{\mathbb{W}_0^r(W_D^{1,q}(\Omega), W^{-1,q}(\Omega))} \le L_R \|u_1 - u_2\|_{L^r(0,T;L^2(\Omega))}.$$
 (2.6)

THEOREM 2.5. [23, Lemma 3.3, Theorem 3.4] The solution operator  $S : L^r(0, T; L^2(\Omega)) \rightarrow W_0^r(W_D^{1,q}(\Omega), W^{-1,q}(\Omega))$  is directionally differentiable and its directional derivative  $\eta = S'(u;h)$  at  $u \in L^r(0,T; L^2(\Omega))$  in direction  $h \in L^r(0,T; L^2(\Omega))$  is given by the unique solution of

$$\dot{\eta}(t) + A\eta(t) + f'(y(t);\eta(t)) = B h(t) \quad a.e. \ in \ (0,T)$$
  
$$\eta(0) = 0, \tag{2.7}$$

with y = S(u). The solution operator of (2.7), namely  $S'(u; \cdot) : L^r(0, T; L^2(\Omega)) \ni h \mapsto \eta \in \mathbb{W}^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega))$  is globally Lipschitz continuous.

Some remarks concerning Assumption 2.1 are in order:

REMARK 2.6. Note that Assumption 2.1.1 does not allow us to consider the  $L^2(Q)$ -Hilbert space for the control, since r > 2 even in two dimensions. We deliberately choose to work with such a setting, although additional assumptions on the nonlinearity f would enable us to set r = 2, cf. [23, Remark 6.5]. We proceed in this way due to the following reason. The condition (2.1) guarantees that the control-tostate mapping is (locally) Lipschitz continuous with range in  $L^{\infty}(Q)$  (see the proof of [23, Lemma 6.4]). This will be indispensable for the derivation of SSC: given a point y in a neighborhood of the state  $\bar{y}$ , we have to be able to make assertions about the distance between y(t, x) and  $\bar{y}(t, x)$  (in the "a.e." sense), see proofs of Lemmas 4.9 and 4.12 below. Further, we observe that r becomes large if q approaches n. We could weaken Assumption 2.1.1 as in [23, Remark 6.5] if e.g.  $RgB = L^2(\Omega)$ . However, one still obtains r > 2, so that the  $L^2(Q)$ -Hilbert space setting for the control is excluded. As mentioned above, extra conditions on f allow for r = 2. Let us point out that in this case one cannot expect  $L^{\infty}(Q)$ -regularity for the state, see [31, Chp. 5].

In the elliptic case, the local Lipschitz continuity of the control-to-state operator with range in the space of essentially bounded functions is crucial too. Thanks to the Stampacchia method, see [31], one can choose  $L^2(\Omega)$  as space for the control to guarantee this. Let us emphasize that the entire second-order analysis for the elliptic version of (P) can be performed in the same way as in the parabolic setting.

REMARK 2.7. In two dimensions, Assumptions 2.1.2-3 are satisfied by the operator  $A = -\operatorname{div} \kappa \nabla$  defined as

$$A: W_D^{1,q}(\Omega) \ni y \mapsto \int_{\Omega} \kappa \nabla y \nabla \cdot \, dx \in W^{-1,q}(\Omega),$$

if  $\Omega \cup \Gamma_N$  is regular in the sense of Gröger, cf. [15], and the coefficient function  $\kappa \in L^{\infty}(\Omega; \mathbb{R}^{n \times n})$  is uniformly elliptic and symmetric. The papers [20, Appendix] and [12] provide many settings such that  $-\operatorname{div} \kappa \nabla$  fulfills Assumptions 2.1.2-3 in three dimensions too, e.g., if  $\Gamma_N = \emptyset$ ,  $\kappa$  is uniformly continuous and may jump across a  $C^1$ -interface, and  $\Omega$  is a strong Lipschitz domain in the sense of [22, Chp. 1.1.9]. For more details, see [23, Remark 6.3]. Note further that A satisfies maximal parabolic  $L^s(0,T; W^{-1,q}(\Omega))$ -regularity for every  $s \in (1,\infty)$ , cf. [13].

REMARK 2.8. The monotony property of f in Assumption 2.1.4 can be replaced by the more general [23, Assumption 2.5], see the proof of [23, Lemma 6.6]. Alternatively, one could require that f satisfies certain growth conditions, cf. [23, Remark 2.6].

REMARK 2.9. Semilinear parabolic PDEs with non-smooth nonlinearities f of the type (2.5) arise for instance in the modeling of combustion processes, see e.g. [33]. In this case, the combustion nonlinearity f features a so-called ignition-temperature  $\Theta > 0$ . For example, f could be identically zero on the interval  $(-\infty, \Theta]$ , i.e., there is no reaction below the ignition-temperature. Once  $\Theta$  is reached, ignition and combustion suddenly occur (so that  $\Theta$  is a kink point of f). This leads to an abrupt change in the physical regime modeled by the non-smooth function f. For more details, we refer to [33, Chp. 1.2]. In Figure 5.1 below we depict an ignition-type nonlinearity with three ignition temperatures.

3. Bouligand-differentiability of the control-to-state operator. This section addresses an essential property of the operator S that will be needed in the proof of the main result, namely its Bouligand-differentiability. To show the latter, we need the following

PROPOSITION 3.1. The function f is Bouligand-differentiable from  $L^{\infty}(Q)$  to  $L^{\beta}(Q)$  for every  $\beta < \infty$ , i.e.,

$$\|f(y+h) - f(y) - f'(y;h)\|_{L^{\beta}(Q)} = o(\|h\|_{L^{\infty}(Q)}) \quad \forall y \in L^{\infty}(Q).$$

In particular, f is directionally differentiable from  $L^{\infty}(Q)$  to  $L^{\beta}(Q)$  for every  $\beta < \infty$ . Furthermore,  $f'(y;h) \in L^{\infty}(Q) \ \forall y, h \in L^{\infty}(Q)$ .

*Proof.* From Assumptions 2.1.4 and 2.1.5 it follows that  $f : \mathbb{R} \to \mathbb{R}$  is Bouligand-differentiable at any  $z \in \mathbb{R}$ , i.e.,

$$\frac{|f(z+v_n) - f(z) - f'(z;v_n)|}{|v_n|} \to 0 \quad \text{as } v_n \to 0.$$
(3.1)

Note that, for all M > 0, it holds

$$|f'(z;v)| \le L_{M+1} |v| \quad \forall v \in \mathbb{R}, \ z \in [-M, M],$$
(3.2)

where  $L_{M+1} > 0$  is given by Assumption 2.1.4. This is due to the definition of the directional derivative and its positive homogeneity w.r.t. direction, see also the proof of [23, Lemma 3.1]. Let now  $\{h_n\} \subset L^{\infty}(Q)$  be an arbitrary sequence with  $h_n \to 0$  in  $L^{\infty}(Q)$  as  $n \to \infty$ . Hence,  $h_n(t, x) \to 0$  a.e. in Q. As a result of (3.1), we have the following convergence

$$g_n(t,x) := \frac{|f(y(t,x) + h_n(t,x)) - f(y(t,x)) - f'(y(t,x);h_n(t,x))|}{\|h_n\|_{L^{\infty}(Q)}} \\ \leq \frac{|f(y(t,x) + h_n(t,x)) - f(y(t,x)) - f'(y(t,x);h_n(t,x))|}{|h_n(t,x)|} \to 0 \quad \text{as } n \to \infty$$

a.e. in Q. Moreover, for n large enough, we have  $|g_n(t,x)| \leq 2L_{\|y\|_{L^{\infty}(Q)}+1}$  a.e. in Q, in view of Assumption 2.1.4 and (3.2). Thus, by Lebesgue's dominated convergence theorem, we obtain

$$g_n \to 0$$
 in  $L^{\beta}(Q)$  as  $n \to \infty$ ,

for any  $\beta < \infty$ , which is the desired assertion. Note that  $f'(y;h) \in L^{\infty}(Q)$  for all  $y, h \in L^{\infty}(Q)$ , in light of (3.2).  $\Box$ 

THEOREM 3.2. The control-to-state mapping  $S: L^r(0,T;L^2(\Omega)) \to \mathbb{W}^r_0(W^{1,q}_D(\Omega),W^{-1,q}(\Omega))$  is Bouligand-differentiable, i.e.,

$$\|S(u+h) - S(u) - S'(u;h)\|_{\mathbb{W}^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega))} = o(\|h\|_{L^r(0,T;L^2(\Omega))})$$

for all  $u \in L^r(0,T;L^2(\Omega))$ .

*Proof.* Let  $u, h \in L^r(0, T; L^2(\Omega))$  be arbitrary, but fixed and set  $y := S(u), y^h := S(u+h)$  and  $\eta := S'(u;h)$ . By subtracting (2.5) and (2.7) from (2.5) with right-hand side u + h we have

$$\frac{d}{dt}(y^h - y - \eta) + A(y^h - y - \eta) = -f(y^h) + f(y) + f'(y;\eta),$$

$$(y^h - y - \eta)(0) = 0.$$
(3.3)

The associated integral equation reads

$$(y^{h} - y - \eta)(t) = \int_{0}^{t} e^{-(t-s)A} \left( -f(y^{h}(s)) + f(y(s)) + f'(y(s);\eta(s)) \right) ds$$

see e.g. [17]. Consequently, one obtains

$$\begin{aligned} \|(y^{h} - y - \eta)(t)\|_{L^{\infty}(\Omega)} & \leq \int_{0}^{t} \|e^{-(t-s)A}\|_{\mathcal{L}(W^{-1,q}(\Omega),L^{\infty}(\Omega))} \Big(\underbrace{\|f(y^{h}(s)) - f(y(s) + \eta(s))\|_{W^{-1,q}(\Omega)}}_{&+ \underbrace{\|f(y(s) + \eta(s)) - f(y(s)) - f'(y(s);\eta(s))\|_{W^{-1,q}(\Omega)}}_{&=:\hat{B}_{h}(s)} \Big) ds \end{aligned}$$

$$(3.4)$$

for all  $t \in [0,T]$ . We assume that  $||h||_{L^r(0,T;L^2(\Omega))} \leq 1$ , since  $h \to 0$  later anyway. Proceeding exactly as in the proof of [23, Lemma 6.6], we find

$$\|y^{h}\|_{C([0,T];L^{\infty}(\Omega))} \leq C(1 + \|u + h\|_{L^{r}(0,T;L^{2}(\Omega))}) \leq C(2 + \|u\|_{L^{r}(0,T;L^{2}(\Omega))}) =: \rho_{1},$$
  
$$\|y + \eta\|_{C([0,T];L^{\infty}(\Omega))} \leq C(1 + \|u\|_{L^{r}(0,T;L^{2}(\Omega))} + L_{u}) =: \rho_{2},$$

where for the last estimate we used Theorem 2.5 combined with (2.4). Now, applying (2.2) with  $M := \max\{\rho_1, \rho_2\}$  yields

$$\hat{A}_{h}(t) \le L_{M} \| (y^{h} - y - \eta)(t) \|_{L^{\infty}(\Omega)} \quad \forall t \in [0, T],$$
(3.5)

in view of  $L^{\infty}(\Omega) \hookrightarrow W^{-1,q}(\Omega)$ . To estimate  $\hat{B}_h$ , let us first fix  $\zeta > 1$  so that  $W_D^{1,q'}(\Omega) \hookrightarrow L^{\zeta}(\Omega)$ , say  $\zeta = 3/2$  and  $\zeta = 2$  in the three and two dimensional case, respectively. Then,  $L^{\zeta'}(\Omega) \hookrightarrow W^{-1,q}(\Omega)$  follows, and by employing the Lipschitz continuity of  $S'(u; \cdot) : L^r(0, T; L^2(\Omega)) \to \mathbb{W}_0^r(W_D^{1,q}(\Omega), W^{-1,q}(\Omega))$ , see Theorem 2.5, we have

$$\begin{aligned} \|\hat{B}_{h}\|_{L^{r}(0,T)} &\leq C \frac{\|f(y+\eta) - f(y) - f'(y;\eta)\|_{L^{r}(0,T;L^{\zeta'}(\Omega))}}{\|\eta\|_{L^{\infty}(Q)}} \|\eta\|_{L^{\infty}(Q)} \\ &\leq C \frac{\|f(y+\eta) - f(y) - f'(y;\eta)\|_{L^{\max\{r,\zeta'\}}(Q)}}{\|\eta\|_{L^{\infty}(Q)}} L_{u} \|h\|_{L^{r}(0,T;L^{2}(\Omega))}, \end{aligned}$$

provided that  $\eta \neq 0$ . Note that for the second inequality we used again (2.4). In case that  $\eta = 0$ , we deduce  $\hat{B}_h = 0$ , by the definition of  $\hat{B}_h$ . Further, we observe that  $\|h\|_{L^r(0,T;L^2(\Omega))} \to 0$  implies  $S'(u;h) \to 0$  in  $L^{\infty}(Q)$ , as a consequence of Theorem 2.5 combined with (2.4). Thanks to  $\zeta > 1$ , we have  $\max\{r, \zeta'\} < \infty$ , and Proposition 3.1 gives in turn

$$\frac{\|\hat{B}_{h}\|_{L^{r}(0,T)}}{\|h\|_{L^{r}(0,T;L^{2}(\Omega))}} \leq C \frac{\|f(y+\eta) - f(y) - f'(y;\eta)\|_{L^{\max\{r,\zeta'\}}(Q)}}{\|\eta\|_{L^{\infty}(Q)}} \to 0,$$
(3.6)

if  $\|h\|_{L^r(0,T;L^2(\Omega))} \to 0$ . Now we return to (3.4), where inserting (3.5) results in

$$\begin{aligned} \|(y^{h} - y - \eta)(t)\|_{L^{\infty}(\Omega)} \\ &\leq L_{M} \int_{0}^{t} \|e^{-(t-s)A}\|_{\mathcal{L}(W^{-1,q}(\Omega),L^{\infty}(\Omega))}\|(y^{h} - y - \eta)(s)\|_{L^{\infty}(\Omega)} \, ds \\ &+ \int_{0}^{t} \|e^{-(t-s)A}\|_{\mathcal{L}(W^{-1,q}(\Omega),L^{\infty}(\Omega))}\hat{B}_{h}(s) \, ds \\ &\leq L_{M} \int_{0}^{t} c \, (t-s)^{-\theta}\|(y^{h} - y - \eta)(s)\|_{L^{\infty}(\Omega)} \, ds \\ &+ \|e^{-A}\|_{L^{r'}(0,T;\mathcal{L}(W^{-1,q}(\Omega),L^{\infty}(\Omega)))}\|\hat{B}_{h}\|_{L^{r}(0,T)} \quad \forall t \in [0,T]. \end{aligned}$$

Here,  $\theta \in (0,1)$  denotes the exponent in (2.3). In view of the latter, the mapping  $t \mapsto e^{-tA}$  belongs indeed to  $L^{r'}(0,T; \mathcal{L}(W^{-1,q}(\Omega), L^{\infty}(\Omega)))$ , since  $r'\theta < 1$ . By means of a generalized Gronwall's inequality, cf. [17, Lemma 7.1.1, p. 188], we have

$$||(y^h - y - \eta)(t)||_{L^{\infty}(\Omega)} \le C ||\hat{B}_h||_{L^r(0,T)} \quad \forall t \in [0,T].$$

Then, by (3.6), the estimate

,

$$\frac{\|y^{h} - y - \eta\|_{C([0,T];L^{\infty}(\Omega))}}{\|h\|_{L^{r}(0,T;L^{2}(\Omega))}} \le C \frac{\|B_{h}\|_{L^{r}(0,T)}}{\|h\|_{L^{r}(0,T;L^{2}(\Omega))}} \to 0 \quad \text{if } \|h\|_{L^{r}(0,T;L^{2}(\Omega))} \to 0 \quad (3.7)$$

follows. Since A satisfies maximal parabolic  $L^r(0,T;W^{-1,q}(\Omega))$ -regularity, see Assumption 2.1.3, we finally arrive at

$$\begin{split} \|y^{h} - y - \eta\|_{\mathbb{W}_{0}^{r}(W_{D}^{1,q}(\Omega),W^{-1,q}(\Omega))} \\ &\leq \|(\partial_{t} + A)^{-1}\|_{\mathcal{L}(L^{r}(0,T;W^{-1,q}(\Omega)),\mathbb{W}_{0}^{r}(W_{D}^{1,q}(\Omega),W^{-1,q}(\Omega)))}\| - f(y^{h}) + f(y) + f'(y;\eta)\|_{L^{r}(0,T;W^{-1,q}(\Omega))} \\ &\leq c \left(\|\hat{A}_{h}\|_{L^{r}(0,T)} + \|\hat{B}_{h}\|_{L^{r}(0,T)}\right) \\ &\leq c \left(\|y^{h} - y - \eta\|_{C([0,T];L^{\infty}(\Omega))} + \|\hat{B}_{h}\|_{L^{r}(0,T)}\right) \\ &= o(\|h\|_{L^{r}(0,T;L^{2}(\Omega))}), \end{split}$$

where we used (3.3), (3.5), (3.7), and (3.6). The proof is now complete.

4. Second-order sufficient optimality conditions. This section is devoted to establishing second-order sufficient conditions (SSC) which guarantee local optimality for (P). Recall that this reads as follows:

$$\begin{array}{ccc}
\min_{u \in L^{r}(0,T;L^{2}(\Omega))} & J(y,u) \\
\text{s.t.} & \dot{y}(t) + A y(t) + f(y(t)) = B u(t) \quad \text{a.e. in } (0,T) \\
& y(0) = 0.
\end{array}$$
(P)

The upcoming second-order analysis relies on the Bouligand-differentiability and the (local) Lipschitz continuity of S with range in  $L^{\infty}(Q)$ . The main results are stated in Theorems 4.15 and 4.18 below. Similarly to [3], we present two versions of second-order sufficient optimality conditions. The first set of conditions involves the positive-definiteness of the Hessian of a "Lagrangian" on the cone of critical directions and applies only to objectives with a particular structure, see Assumption 4.13 below. The second set of SSC allows for general (smooth) objectives. In this case, the Hessian of the "Lagrangian" is supposed to be coercive on a larger cone, cf. Assumption 4.16 below. Moreover, Theorems 4.15 and 4.18 below are based on the strong stationarity result from [23] and the premise that the adjoint state fulfills a suitable sign condition (on a certain subset of Q). Let us point out that if f is twice continuously differentiable, then both versions of SSC derived in this section coincide with the classical SSC, see Remark 4.19 below.

We begin by recalling the necessary optimality condition for (P) (in form of strong stationarity) established in [23].

THEOREM 4.1. [23, Thm. 5.3, Thm. 6.7] Suppose that the range of B is dense in  $W^{-1,q}(\Omega)$ . Let  $\bar{u} \in L^r(0,T; L^2(\Omega))$  be locally optimal for (P) with associated state  $\bar{y} = S(\bar{u}) \in W^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega))$ . Then there exists a unique adjoint state  $p \in W^{r'}_T(W^{1,q'}_D(\Omega), W^{-1,q'}(\Omega))$  and a unique multiplier  $\lambda \in L^{r'}(0,T; L^s(\Omega))$  with  $s = \frac{nq}{nq-n-q}$  such that

$$\dot{\bar{y}} + A\,\bar{y} + f(\bar{y}) = B\,\bar{u}, \quad \bar{y}(0) = 0,$$
(4.1a)

$$-\dot{p} + A^* p + \lambda = \partial_y J(\bar{y}, \bar{u}), \quad p(T) = 0, \tag{4.1b}$$

$$\lambda(t,x) \in [f'_+(\bar{y}(t,x))\,p(t,x), f'_-(\bar{y}(t,x))\,p(t,x)] \quad a.e. \ in \ Q, \tag{4.1c}$$

$$B^* p + \partial_u J(\bar{y}, \bar{u}) = 0, \tag{4.1d}$$

where, for an arbitrary  $z \in \mathbb{R}$ , the right- and left-sided derivative of  $f : \mathbb{R} \to \mathbb{R}$  are defined through  $f'_+(z) := f'(z; 1)$  and  $f'_-(z) := -f'(z; -1)$ , respectively.

In the remaining of the section, let  $(\bar{u}, \bar{y}, \lambda, p)$  be a fixed point which satisfies the system (4.1) and possesses the same regularity as in Theorem 4.1.

LEMMA 4.2. For any  $u \in L^r(0,T;L^2(\Omega))$ , there exists  $\gamma \in [0,1]$  such that

$$J(y,u) - J(\bar{y},\bar{u}) \ge \int_{Q} p(t,x) \left( f(\bar{y}) - f(y) + f'(\bar{y};y-\bar{y}) \right)(t,x) \ d(t,x) + \frac{1}{2} J''(y_{\gamma},u_{\gamma})(y-\bar{y},u-\bar{u})^{2},$$

$$(4.2)$$

where we abbreviate y := S(u) and  $(y_{\gamma}, u_{\gamma}) := (\bar{y}, \bar{u}) + \gamma((y, u) - (\bar{y}, \bar{u}))$ . *Proof.* Let  $u \in L^r(0, T; L^2(\Omega))$  be arbitrary, but fixed. Using the optimality system, we find

$$\begin{aligned} \partial_{y}J(\bar{y},\bar{u})(y-\bar{y}) &+ \partial_{u}J(\bar{y},\bar{u})(u-\bar{u}) \\ &= \langle -\dot{p} + A^{*}p + \lambda, y-\bar{y} \rangle_{L^{r}(0,T;W_{D}^{1,q}(\Omega))} - \langle B^{*}p, u-\bar{u} \rangle_{L^{r}(0,T;L^{2}(\Omega))} \\ &= \langle p, \dot{y} - \dot{\bar{y}} + A(y-\bar{y}) - B(u-\bar{u}) \rangle_{L^{r}(0,T;W^{-1,q}(\Omega))} + \langle \lambda, y-\bar{y} \rangle_{L^{r}(0,T;W_{D}^{1,q}(\Omega))} \\ &\geq \langle p, f(\bar{y}) - f(y) \rangle_{L^{r}(0,T;W^{-1,q}(\Omega))} + \int_{Q} p(t,x) f'(\bar{y};y-\bar{y})(t,x) \ d(t,x). \end{aligned}$$

For the last equality, we applied the formula of integration by parts from [2, Proposition 5.1] in combination with the initial and final time conditions in (2.5) and (4.1b), respectively. The above inequality can be deduced from the state equation (2.5) and (4.1c) together with the positive homogeneity of the directional derivative. The desired assertion follows now from the continuous Fréchet-differentiability of J, cf. Assumption 2.1.7.  $\Box$ 

The key idea in the proofs of Theorems 4.15 and 4.18 below is to write the integral in (4.2) as the sum of a nonnegative term,  $-1/2 \int_{\mathcal{M}} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};u-\bar{u})(t,x)^2 d(t,x)$ , and  $o(||u-\bar{u}||^2_{L^r(0,T;L^2(\Omega))})$ , where  $\mathcal{M}$  is a suitable subset of Q. In preparation therefor, we discuss the first term on the right-hand side in (4.2) on three different subsets of Q, by mainly distinguishing between those (t,x) for which f is differentiable or not at  $\bar{y}(t,x)$ , see (4.3), (4.7), and (4.10) below.

We identify the non-smooth and smooth points of the function f by means of the following sets:

$$\mathcal{N} := \{ z \in \mathbb{R} \mid f \text{ is not differentiable at } z \}, \\ \mathcal{S} := \{ z \in \mathbb{R} \mid f \text{ is differentiable at } z \}.$$

Recall that  $\mathcal{N}$  is at most countable, cf. Assumption 2.1.5. This ensures that the sets in (4.3), (4.7), and (4.10) below are measurable.

Next, we introduce the notion of local convexity/concavity for the nonlinearity f. This will play a crucial role in the next two lemmas.

DEFINITION 4.3. We say that the function f is convex around  $y \in \mathbb{R}$  if there exists  $\rho > 0$  so that f is convex on the interval  $(y - \rho, y + \rho)$ . Analogously, we say that f is concave around  $y \in \mathbb{R}$  if there exists  $\rho > 0$  so that f is concave on the interval  $(y - \rho, y + \rho)$ .

We also define (up to sets of measure zero) the following subset of Q:

$$Q_{\mathfrak{n}} := \{(t, x) \in Q \mid \exists z \in \mathcal{N} \text{ so that } \bar{y}(t, x) = z\}.$$
(4.3)

LEMMA 4.4. Suppose that at any  $z \in \mathcal{N}$ , the function f is either convex or concave around z with radius  $\rho_z > 0$ . If  $\inf_{z \in \mathcal{N}} \rho_z > 0$ , then there exists  $\varepsilon > 0$  so that, for any  $y \in L^{\infty}(Q)$  with  $\|y - \bar{y}\|_{L^{\infty}(Q)} < \varepsilon$ , there holds

$$p(t,x)(f(\bar{y}) - f(y) + f'(\bar{y};y - \bar{y}))(t,x) \ge 0$$
 a.e. in  $Q_n$ 

Proof. Let  $z \in \mathcal{N}$  be a non-smooth point of f for which the set  $\mathfrak{m}_z := \{(t,x) \in Q \mid \overline{y}(t,x) = z\}$  has positive measure. If f is convex around z, then straightforward computation shows that  $f'_+(z) > f'_-(z)$ , since z is a non-differentiable point. Thus,  $f'_+(\overline{y}(t,x)) > f'_-(\overline{y}(t,x))$  a.e. in  $\mathfrak{m}_z$ . On the other hand, from (4.1c) we deduce that the interval  $[f'_+(\overline{y}(t,x)) p(t,x), f'_-(\overline{y}(t,x)) p(t,x)]$  is nonempty for a.a.  $(t,x) \in Q$ . Thus,

$$p(t,x) \le 0 \quad \text{a.e. in } \mathfrak{m}_z. \tag{4.4}$$

Since f is convex on  $(z - \rho_z, z + \rho_z)$ , it holds

$$f(v) - f(z) \ge f'(z; v - z)$$
 for all  $v \in (z - \rho_z, z + \rho_z)$ . (4.5)

Now, we define  $\varepsilon := \inf_{z \in \mathcal{N}} \rho_z > 0$ . Let  $y \in L^{\infty}(Q)$  with  $||y - \bar{y}||_{L^{\infty}(Q)} < \varepsilon$  be arbitrary, but fixed. Then,  $y(t, x) \in (z - \rho_z, z + \rho_z)$  a.e. in  $\mathfrak{m}_z$ , and from (4.5) combined with (4.4) we deduce

$$p(t,x) \big( f(\bar{y}) - f(y) + f'(\bar{y}; y - \bar{y}) \big)(t,x) \ge 0 \quad \text{a.e. in } \mathfrak{m}_z.$$
(4.6)

Analogously, if f is concave around z, one has  $p(t, x) \ge 0$  a.e. in  $\mathfrak{m}_z$ , and (4.6) follows in the same way as above. Since  $Q_{\mathfrak{n}} = \bigcup_{z \in \mathcal{N}} \mathfrak{m}_z$  by definition, the proof is now complete.  $\Box$ 

REMARK 4.5. In view of (4.4), one always knows the sign of the adjoint state p a.e. where  $\bar{y}(t,x)$  is a non-differentiable point of f. Unfortunately, this is no longer the case in the smooth points. No matter how close  $\bar{y}(t,x) \in S$  is to some  $z \in N$ , a sign for p(t,x) could not be provided. This turns out to be a problem precisely for those (t,x) for which  $\bar{y}(t,x)$  is "too close" to the non-smooth points, as explained in Remark 4.8.(iii) below. This is why we need to assume a certain sign condition for p on this critical subset of Q, below also known as  $Q_{s,\delta}$ .

ASSUMPTION 4.6. Suppose that, at any  $z \in \mathcal{N}$ , the function f is either convex or concave around z with radius  $\rho_z > 0$ . Moreover, assume that, for any  $z \in \mathcal{N}$ , there exists  $\delta_z \in (0, \rho_z)$  so that the following conditions are fulfilled:

1.  $\inf_{z \in \mathcal{N}} \rho_z - \delta_z > 0$ ,

2. if f is convex around z:  $p(t, x) \leq 0$  a.e. where  $\bar{y}(t, x) \in (z - \delta_z, z + \delta_z) \setminus \{z\}$ , 3. if f is concave around z:  $p(t, x) \geq 0$  a.e. where  $\bar{y}(t, x) \in (z - \delta_z, z + \delta_z) \setminus \{z\}$ .

REMARK 4.7. Given three successive points in  $\mathcal{N}$ , say  $z_0$ ,  $z_1$ , and  $z_2$ , the value  $\rho_{z_1} > 0$  should be chosen in all what follows so that  $z_0 < z_1 - \rho_{z_1} < z_1 < z_1 + \rho_{z_1} < z_2$  holds.

From now on,  $\{\delta_z\}_{z \in \mathcal{N}}$  is supposed to be a fixed set of strict positive, as small as possible values that satisfy Assumption 4.6. By means thereof, we define (up to sets of measure zero):

$$Q_{\mathfrak{s},\delta} := \{(t,x) \in Q \mid \bar{y}(t,x) \in \bigcup_{z \in \mathcal{N}} (z - \delta_z, z + \delta_z) \setminus \{z\}\}.$$
(4.7)

REMARK 4.8. (i) For each  $z \in \mathcal{N}$ , it is desirable to choose  $\delta_z > 0$  so (small) that  $\{(t,x) \in Q \mid \overline{y}(t,x) \in (z-\delta_z, z+\delta_z) \setminus \{z\}\}$  has measure zero, if possible. In this case, Assumption 4.6.2-3 is automatically fulfilled at  $z \in \mathcal{N}$ .

(ii) The assumption on the sign of p corresponds to an assumption from [21], which was made in the context of deriving SSC for the elliptic obstacle problem. We refer here to [21, Assumption 1.(iii)] and the proof of [21, (2.26)]. In Section 5 below we will provide settings for which Assumption 4.6.2-3 is guaranteed.

(iii) On the set  $Q \setminus (Q_n \cup Q_{\mathfrak{s},\delta})$  we do not need any (sign) conditions. Here we can evaluate the term  $(f(\bar{y}) - f(y) + f'(\bar{y}; y - \bar{y})(t, x))$  by means of a Taylor expansion for any  $y \in B_{L^{\infty}(Q)}(\bar{y}, \varepsilon)$ , where  $\varepsilon > 0$  is chosen appropriately, see (4.13) in the proof of Lemma 4.12 below. Unfortunately, this cannot be done on the critical set  $Q_{\mathfrak{s},\delta}$ : it is not clear if an  $L^{\infty}(Q)$ -neighborhood of  $\bar{y}$  exists so that, for all y in this neighborhood, it holds  $[\bar{y}(t, x), y(t, x)] \subset S$  a.e. in  $Q_{\mathfrak{s},\delta}$ .

(iv) Let us point out that if there exists  $\delta > 0$  so that  $|\bar{y}(t,x) - z| \geq \delta$  a.e. in  $Q \setminus Q_n$ for all  $z \in \mathcal{N}$ , then Assumption 4.6 is no longer needed, and Lemma 4.9 below can be omitted. In this case, we can find a neighborhood as depicted above. Then we can argue as in the proof of Lemma 4.12 below and obtain the therein showed result for the entire set  $Q \setminus Q_n$ .

LEMMA 4.9. Let Assumption 4.6 hold true. Then, there exists  $\varepsilon > 0$  so that, for all  $y \in L^{\infty}(Q)$  with  $\|y - \bar{y}\|_{L^{\infty}(Q)} < \varepsilon$ , it holds

$$p(t,x)\big(f(\bar{y}) - f(y) + f'(\bar{y};y-\bar{y})\big)(t,x) \ge 0 \quad a.e. \text{ in } Q_{\mathfrak{s},\delta}$$

*Proof.* We define  $\varepsilon := \inf_{z \in \mathcal{N}} \rho_z - \delta_z > 0$  and consider  $y \in L^{\infty}(Q)$  with  $\|y - \bar{y}\|_{L^{\infty}(Q)} < \varepsilon$  arbitrary, but fixed. Let  $z \in \mathcal{N}$  and denote by  $\mathfrak{q}_z$  the set  $\{(t, x) \in Q \mid \bar{y}(t, x) \in (z - \delta_z, z + \delta_z) \setminus \{z\}\}$  (up to sets of measure zero). Then, due to  $|y(t, x) - z| - |\bar{y}(t, x) - z| \leq \|y - \bar{y}\|_{L^{\infty}(Q)} < \rho_z - \delta_z$  a.e. in Q, we have

$$|y(t,x) - z| < \rho_z \quad \text{a.e. in } \mathfrak{q}_z. \tag{4.8}$$

If f is convex around z, the inequality  $f(v) - f(w) \ge f'(w; v - w)$  is true for all  $v, w \in (z - \rho_z, z + \rho_z)$ . Since  $\rho_z > \delta_z$  by assumption, the definition of  $\mathfrak{q}_z$  together with (4.8) and the sign assumption on p, cf. Assumption 4.6.2, now yield

$$p(t,x)(f(\bar{y}) - f(y) + f'(\bar{y}; y - \bar{y}))(t,x) \ge 0 \quad \text{a.e. in } q_z.$$
(4.9)

In case that f is concave around z, one arrives at (4.9) in the same way as in the convex case, by making use of Assumption 4.6.3. Since  $z \in \mathcal{N}$  was arbitrary and  $Q_{\mathfrak{s},\delta} = \bigcup_{z \in \mathcal{N}} \mathfrak{q}_z$ , the desired assertion follows now from (4.9).  $\Box$ 

Given the set of strict positive values  $\{\delta_z\}_{z \in \mathcal{N}}$  from Assumption 4.6, we define (up to sets of measure zero):

$$Q_{\mathfrak{s}} := \{(t, x) \in Q \mid |\bar{y}(t, x) - z| \ge \delta_z \ \forall z \in \mathcal{N}\}.$$

$$(4.10)$$

ASSUMPTION 4.10. From now on, we assume that, in addition to Assumption 4.6,  $\hat{\varepsilon} := \inf_{z \in \mathcal{N}} \delta_z/2 > 0$  holds, and that the nonlinearity f is twice continuously differentiable on  $\{v \in \mathbb{R} | |v - z| \ge \delta_z/2 \quad \forall z \in \mathcal{N}\}$ . Furthermore, suppose that there exists L > 0 such that

$$|f'(w) - f'(\bar{y}(t,x))| \le L |w - \bar{y}(t,x)| \quad \forall w \in (\bar{y}(t,x) - \hat{\varepsilon}, \bar{y}(t,x) + \hat{\varepsilon})$$
(4.11)  
11

f.a.a.  $(t, x) \in Q_{\mathfrak{s}}$ .

As a direct consequence of (4.11), we have

$$|f''(\bar{y}(t,x))| = \lim_{\tau \to 0} \frac{|f'(\bar{y}(t,x) + \tau) - f'(\bar{y}(t,x))|}{|\tau|} \le L \quad \text{f.a.a.} \ (t,x) \in Q_{\mathfrak{s}}.$$
 (4.12)

Hence, the mapping  $(t, x) \in Q_{\mathfrak{s}} \mapsto f''(\bar{y}(t, x)) \in \mathbb{R}$  belongs to  $L^{\infty}(Q_{\mathfrak{s}})$ .

REMARK 4.11. We point out that the inequality (4.11) is satisfied if  $\mathcal{N} \cap (-\|\bar{y}\|_{L^{\infty}(Q)}, \|\bar{y}\|_{L^{\infty}(Q)})$  is finite. To see this, we define

$$D := \left[ -\|\bar{y}\|_{L^{\infty}(Q)}, z_{1} - \frac{\delta_{1}}{2} \right] \cup \left[ z_{1} + \frac{\delta_{1}}{2}, z_{2} - \frac{\delta_{2}}{2} \right] \cup \dots \left[ z_{i} + \frac{\delta_{i}}{2}, z_{i+1} - \frac{\delta_{i+1}}{2} \right]$$
$$\dots \cup \left[ z_{m-1} + \frac{\delta_{m-1}}{2}, z_{m} - \frac{\delta_{m}}{2} \right] \cup \left[ z_{m} + \frac{\delta_{m}}{2}, \|\bar{y}\|_{L^{\infty}(Q)} \right],$$

where  $z_i$  is the *i*-th non-smooth point in  $\mathcal{N} \cap (-\|\bar{y}\|_{L^{\infty}(Q)}, \|\bar{y}\|_{L^{\infty}(Q)})$  (in increasing order) and  $\delta_i := \delta_{z_i}$ . Since f'' is continuous on D by assumption, f' is Lipschitz continuous on each of the intervals in D, and (4.11) follows from  $Q_{\mathfrak{s}} \subset \{(t,x) \in Q \mid \bar{y}(t,x) \in D\}$  and the definition of  $\hat{\varepsilon}$ . Note that if  $\mathcal{N}$  is finite, the entire Assumption 4.10 is fulfilled, excepting the continuous differentiability of f.

The next lemma is the last essential step before proving the main result.

LEMMA 4.12. Let Assumption 4.10 be satisfied. Then it holds

$$\begin{split} \int_{Q_s} p(t,x) \big( f(\bar{y}) - f(S(u)) + f'(\bar{y};S(u) - \bar{y}) \big)(t,x) \, d(t,x) \\ &= -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};u - \bar{u})(t,x)^2 \, d(t,x) + o(\|u - \bar{u}\|_{L^r(0,T;L^2(\Omega))}^2). \end{split}$$

*Proof.* Let  $y \in L^{\infty}(Q)$  with  $||y - \bar{y}||_{L^{\infty}(Q)} < \hat{\varepsilon}$  be arbitrary, but fixed. Note that  $\hat{\varepsilon} = \inf_{z \in \mathcal{N}} \delta_z/2 > 0$ , cf. Assumption 4.10. From

$$|\bar{y}(t,x)-z|-|y(t,x)-z| \le \|y-\bar{y}\|_{L^{\infty}(Q)} < \delta_z/2 \quad \text{ a.e. in } Q, \ \forall z \in \mathcal{N}$$

we deduce that  $|y(t,x) - z| > \delta_z/2$  a.e. in  $Q_{\mathfrak{s}}, \forall z \in \mathcal{N}$ . Since for all  $z \in \mathcal{N}$  it holds  $|\bar{y}(t,x) - z| \geq \delta_z$  and  $|y(t,x) - \bar{y}(t,x)| < \delta_z/2$  a.e. in  $Q_{\mathfrak{s}}$ , every point between  $\bar{y}(t,x)$  and y(t,x) belongs to  $\{v \in \mathbb{R} | |v - z| > \delta_z/2 \forall z \in \mathcal{N}\}$  f.a.a.  $(t,x) \in Q_{\mathfrak{s}}$ . Thus, f is twice continuously differentiable on  $[\bar{y}(t,x), y(t,x)]$  f.a.a.  $(t,x) \in Q_{\mathfrak{s}}$ , in view of Assumption 4.10. This allows us to write the Taylor formula

$$f(y(t,x)) = f(\bar{y}(t,x)) + f'(\bar{y}(t,x))(y(t,x) - \bar{y}(t,x)) + 1/2f''(\bar{y}(t,x))(y(t,x) - \bar{y}(t,x))^2 + o((y(t,x) - \bar{y}(t,x))^2) \quad \text{a.e. in } Q_{\mathfrak{s}}.$$
(4.13)

Further, from (2.6) and (2.4) we know that there exists a constant  $C = C(\bar{u}) > 0$  such that

$$|S(v) - \bar{y}||_{L^{\infty}(Q)} \le C ||v - \bar{u}||_{L^{r}(0,T;L^{2}(\Omega))}$$
(4.14)

for all  $v \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u},1)}$ . Let now  $\epsilon := \min\{\hat{\varepsilon}/2C,1\} > 0$  and  $u \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u},\epsilon)}, u \neq \bar{u}$ , be arbitrary, but fixed. Then, due to (4.14), we have  $\|S(u) - \bar{y}\|_{L^{\infty}(Q)} < \hat{\varepsilon}$  and

as a result of (4.13), it holds

$$f(S(u)(t,x)) = f(\bar{y}(t,x)) + f'(\bar{y}(t,x))(S(u)(t,x) - \bar{y}(t,x)) + 1/2f''(\bar{y}(t,x))(S(u)(t,x) - \bar{y}(t,x))^{2} + \underbrace{o((S(u)(t,x) - \bar{y}(t,x))^{2})}_{=:h_{u}(t,x)} \quad \text{a.e. in } Q_{\mathfrak{s}}.$$

$$(4.15)$$

Hence,

$$\int_{Q_s} p(t,x) \left( f(\bar{y}) - f(S(u)) + f'(\bar{y}; S(u) - \bar{y}) \right)(t,x) d(t,x) = -\frac{1}{2} \underbrace{\int_{Q_s} p(t,x) f''(\bar{y}(t,x)) \left( S(u)(t,x) - \bar{y}(t,x) \right)^2 d(t,x)}_{\hat{A}_u} - \underbrace{\int_{Q_s} p(t,x) h_u(t,x) d(t,x)}_{\hat{B}_u} .$$
(4.16)

In the following, we discuss the terms  $\hat{A}_u$  and  $\hat{B}_u$  separately. We begin with the first term on the right-hand side of (4.16). For convenience, we abbreviate

$$F(u) := S(u) - S(\bar{u}) - S'(\bar{u}; u - \bar{u}).$$

In view of Theorem 3.2, it holds

$$||F(u)||_{L^{\infty}(Q)} = o(||u - \bar{u}||_{L^{r}(0,T;L^{2}(\Omega))}),$$
(4.17)

where we again employed (2.4). Now we write  $\hat{A}_u$  as

$$\hat{A}_{u} = \int_{Q_{s}} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};u-\bar{u})(t,x)^{2} d(t,x) \\
+ \underbrace{\int_{Q_{s}} p(t,x) f''(\bar{y}(t,x)) F(u)(t,x) \Big(F(u)(t,x) + 2S'(\bar{u};u-\bar{u})(t,x)\Big) d(t,x) \,.}_{D_{u}} \tag{4.18}$$

In light of the global Lipschitz continuity of  $S'(\bar{u}; \cdot)$ , see Theorem 2.5, in combination with the embedding (2.4), we obtain

$$\frac{|D_u|}{\|u - \bar{u}\|^2} \le \|p\|_{L^1(Q)} \|f''(\bar{y}(\cdot))\|_{L^{\infty}(Q_s)} \frac{\|F(u)\|_{L^{\infty}(Q)}}{\|u - \bar{u}\|} \frac{(\|F(u)\|_{L^{\infty}(Q)} + 2\|S'(\bar{u}; u - \bar{u})\|_{L^{\infty}(Q)})}{\|u - \bar{u}\|} \xrightarrow{\longrightarrow 0 \quad \text{if } \|u - \bar{u}\|_{L^r(0,T;L^2(\Omega))} \to 0,}$$

$$(4.19)$$

as a result of (4.17). Next, we address the term  $\hat{B}_u$ . From (4.15) we read that  $h_u(t,x) = 0$  a.e. where  $S(u)(t,x) = \bar{y}(t,x)$  in  $Q_s$ . F.a.a.  $(t,x) \in Q_s$  for which  $S(u)(t,x) \neq \bar{y}(t,x)$ , we have

$$\frac{|h_u(t,x)|}{\|u-\bar{u}\|^2} \le C^2 \frac{|h_u(t,x)|}{(S(u)(t,x)-\bar{y}(t,x))^2},$$

in view of (4.14). Since  $||u - \bar{u}||_{L^r(0,T;L^2(\Omega))} \to 0$  implies  $S(u)(t,x) \to \bar{y}(t,x)$  a.e. in Q, the definition of  $h_u$  gives

$$h_u(t,x) = o(\|u - \bar{u}\|_{L^r(0,T;L^2(\Omega))}^2) \quad \text{f.a.a.} \ (t,x) \in Q_{\mathfrak{s}}.$$
(4.20)
13

Further, from (4.15) and (4.12) we infer, by applying the mean value theorem, that

$$\begin{aligned} |h_u(t,x)| &\leq |f(S(u)(t,x)) - f(\bar{y}(t,x)) - f'(\bar{y}(t,x))(S(u)(t,x) - \bar{y}(t,x)) \\ &+ L/2(S(u)(t,x) - \bar{y}(t,x))^2 \\ &= |(f'(\tilde{y}_u(t,x)) - f'(\bar{y}(t,x)))(S(u)(t,x) - \bar{y}(t,x))| \\ &+ L/2(S(u)(t,x) - \bar{y}(t,x))^2 \quad \text{a.e. in } Q_{\mathfrak{s}}, \end{aligned}$$

where  $\tilde{y}_u(t,x) := \gamma_u(t,x)(S(u)(t,x) - \bar{y}(t,x)) + \bar{y}(t,x)$ , with some  $\gamma_u(t,x) \in (0,1)$ . Due to  $\|S(u) - \bar{y}\|_{L^{\infty}(Q)} < \hat{\varepsilon}$ , we have

$$|\tilde{y}_u(t,x) - \bar{y}(t,x)| = |\gamma_u(t,x)(S(u)(t,x) - \bar{y}(t,x))| < \hat{\varepsilon} \quad \text{f.a.a.} \ (t,x) \in Q_{\mathfrak{s}},$$

whence, by (4.11),

$$\begin{aligned} |h_u(t,x)| &\leq L|\tilde{y}_u(t,x) - \bar{y}(t,x)| |S(u)(t,x) - \bar{y}(t,x)| + L/2(S(u)(t,x) - \bar{y}(t,x))^2 \\ &\leq 3L/2 \, \|S(u) - \bar{y}\|_{L^{\infty}(Q)}^2 \quad \text{f.a.a.} \ (t,x) \in Q_{\mathfrak{s}} \end{aligned}$$

follows. Now, (4.14) ensures that

$$\frac{|h_u(t,x)|}{\|u-\bar{u}\|^2} \le 3LC^2/2 \quad \text{a.e. in } Q_{\mathfrak{s}}.$$
(4.21)

By means of Lebesgue's dominated convergence theorem, we deduce

$$\frac{\|h_u\|_{L^{\beta}(Q_s)}}{\|u-\bar{u}\|^2} \longrightarrow 0 \quad \text{if } \|u-\bar{u}\|_{L^r(0,T;L^2(\Omega))} \to 0$$
(4.22)

for any  $\beta < \infty$ , in view of (4.20) and (4.21). Since  $r, q < \infty$ , cf. Assumption 2.1.1, we can fix  $\beta < \infty$  such that  $\beta' = \min\{r', q'\} > 1$ . Then,  $p \in L^{r'}(0, T; W_D^{1,q'}(\Omega)) \hookrightarrow L^{\beta'}(Q)$ , by Theorem 4.1. With (4.22) we can thus infer that

$$\frac{|\ddot{B}_{u}|}{\|u-\bar{u}\|^{2}} \le \|p\|_{L^{\beta'}(Q)} \frac{\|h_{u}\|_{L^{\beta}(Q_{\mathfrak{s}})}}{\|u-\bar{u}\|^{2}} \longrightarrow 0 \quad \text{if } \|u-\bar{u}\|_{L^{r}(0,T;L^{2}(\Omega))} \to 0.$$
(4.23)

By inserting (4.18), (4.19) and (4.23) in (4.16), we finally arrive at the desired result.  $\Box$ 

We are now in the position to establish the first version of second-order sufficient optimality conditions for (P). Besides Assumptions 4.6 and 4.10, the SSC contain structural assumptions on J and a positive-definiteness condition for the Hessian of a functional which corresponds to the classical Lagrange-functional, see Remark 4.19 below.

ASSUMPTION 4.13. In addition to Assumptions 4.6 and 4.10, we require that

1. The objective  $J : L^r(0,T; L^{\infty}(\Omega)) \times L^r(0,T; L^2(\Omega)) \to \mathbb{R}$  is given by J(y,u) = g(y) + j(u), where  $g : L^r(0,T; L^{\infty}(\Omega)) \to \mathbb{R}$  and  $j : L^r(0,T; L^2(\Omega)) \to \mathbb{R}$  are both twice continuously Fréchet-differentiable. There exists  $\nu > 0$  with

$$j''(\bar{u})(h,h) \ge \nu \|h\|_{L^{r}(0,T;L^{2}(\Omega))}^{2} \quad \forall h \in L^{r}(0,T;L^{2}(\Omega)).$$
(4.24)

2. For all  $h \in L^r(0,T;L^2(\Omega)) \setminus \{0\}$  and  $\eta = S'(\bar{u};h)$  with  $g'(\bar{y})\eta + j'(\bar{u})h = 0$ , it holds

$$g''(\bar{y})(\eta,\eta) + j''(\bar{u})(h,h) - \int_{Q_s} p(t,x) f''(\bar{y}(t,x))\eta(t,x)^2 d(t,x) > 0, \quad (4.25)$$

where  $Q_{\mathfrak{s}}$  is the set associated with  $\{\delta_z\}_{z\in\mathcal{N}}$  given by (4.10).

REMARK 4.14. (i) The coercivity condition (4.24) in Assumption 4.13.1 is satisfied by the quadratic functional  $\hat{J}(y, u) := \frac{1}{2} ||y - y_d||^2_{L^2(Q)} + \frac{\nu}{2} ||u - u_d||^2_{H^1(0,T;L^2(\Omega))}$ , where  $\nu > 0$ ,  $y_d \in L^2(Q)$ , and  $u_d \in H^1(0,T;L^2(\Omega))$ , due to the embedding  $H^1(0,T;L^2(\Omega)) \hookrightarrow L^{\infty}(0,T;L^2(\Omega))$ . In this case, we need to consider the control space  $H^1(0,T;L^2(\Omega))$  instead of  $L^r(0,T;L^2(\Omega))$ . Note that, for any  $q \in (n,\infty)$ , one can always find rsatisfying (2.1) so that  $H^1(0,T;L^2(\Omega)) \hookrightarrow L^r(0,T;L^2(\Omega))$  holds. Thus, the entire analysis, including the second-order analysis, remains the same with one exception: the result in Theorem 4.15 below is true, provided that we replace (4.27) below by

$$\hat{J}(\bar{y},\bar{u}) + \alpha \|u - \bar{u}\|_{H^1(0,T;L^2(\Omega))}^2 \le \hat{J}(S(u),u) \quad \forall u \in B_{H^1(0,T;L^2(\Omega))}(\bar{u},R).$$
(4.26)

(ii) If B is injective and  $J = \hat{J}$ , it can be shown, by following the lines of the proof of [10, Corollary 4.14], that  $\bar{u}$  is a strong local solution, cf. [10], i.e., there exist  $\alpha' > 0$  and R' > 0 so that

$$\begin{aligned} \hat{J}(\bar{y},\bar{u}) + \alpha' \|u - \bar{u}\|_{H^1(0,T;L^2(\Omega))}^2 &\leq \hat{J}(S(u),u) \quad \forall u \in H^1(0,T;L^2(\Omega)) \\ & \text{with } \|S(u) - \bar{y}\|_{L^2(O)} < R'. \end{aligned}$$

(provided that the corresponding assumptions in Theorem 4.15 below are satisfied). We point out that this assertion is sharper than the quadratic growth condition (4.26) or than those in [3] and [21].

THEOREM 4.15. Let  $(\bar{u}, \bar{y}, \lambda, p)$  satisfy the first-order optimality system (4.1) given by Theorem 4.1. If Assumptions 4.6, 4.10 and 4.13 are fulfilled, then there exist  $\alpha > 0$ and R > 0 such that

$$J(\bar{y},\bar{u}) + \alpha \|u - \bar{u}\|_{L^{r}(0,T;L^{2}(\Omega))}^{2} \leq J(S(u),u) \quad \forall u \in B_{L^{r}(0,T;L^{2}(\Omega))}(\bar{u},R).$$
(4.27)

In particular,  $\bar{u}$  is locally optimal for (P).

*Proof.* The proof is inspired by the proof of [21, Thm. 2.12], see also the proof of [3, Thm. 4.12]. We assume that (4.27) is not satisfied. Thus, there exists a sequence  $\{u_k\}_k \subset L^r(0,T;L^2(\Omega))$  with  $u_k \to \overline{u}$  in  $L^r(0,T;L^2(\Omega))$  and

$$J(\bar{y}, \bar{u}) + \frac{1}{k} \|u_k - \bar{u}\|_{L^r(0,T;L^2(\Omega))}^2 > J(y_k, u_k) \quad \forall k \in \mathbb{N},$$

where  $y_k := S(u_k)$  for the rest of the proof. For simplicity, we define  $\sigma_k := ||u_k - \bar{u}||_{L^r(0,T;L^2(\Omega))}$  and  $h_k := \frac{u_k - \bar{u}}{\sigma_k}$ . Then, the above inequality reads

$$J(\bar{y},\bar{u}) + \frac{1}{k}\sigma_k^2 > J(y_k,u_k) \quad \forall k \in \mathbb{N}.$$
(4.28)

Since  $||h_k||_{L^r(0,T;L^2(\Omega))} = 1$  and  $L^r(0,T;L^2(\Omega))$  is reflexive, see (2.1), we can extract a subsequence, denoted by the same symbol, so that

$$h_k \rightarrow h$$
 in  $L^r(0,T; L^2(\Omega))$  as  $k \rightarrow \infty$ . (4.29)

By arguing in the same way as in the proof of [23, Lemma 2.13], one can prove the weak continuity of the operator  $S'(\bar{u}; \cdot) : L^r(0, T; L^2(\Omega)) \to \mathbb{W}^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega))$ . This implies

$$S'(\bar{u};h_k) \to S'(\bar{u};h) \quad \text{in } C([0,T];L^{\infty}(\Omega)) \quad \text{as } k \to \infty,$$

$$15$$
(4.30)

as a consequence of (4.29) and (2.4). Moreover, according to Theorem 3.2, S is Bouligand-differentiable, and thus,

$$\frac{S(u_k) - S(\bar{u}) - S'(\bar{u}; u_k - \bar{u})}{\sigma_k} \to 0 \quad \text{in } L^{\infty}(Q) \text{ as } k \to \infty$$

Hence, by the positive homogeneity of the directional derivative, we have

$$\frac{S(u_k) - S(\bar{u})}{\sigma_k} \to S'(\bar{u}; h) \quad \text{in } L^{\infty}(Q) \text{ as } k \to \infty.$$
(4.31)

Next, we show that  $(h, S'(\bar{u}; h))$  is a 'critical direction'. As a result of (4.28), it holds

$$g'(\bar{y})(y_k - \bar{y}) + j'(\bar{u})(u_k - \bar{u}) < \frac{\sigma_k^2}{k} - \frac{1}{2}g''(\tilde{y}_k)(y_k - \bar{y})^2 - \frac{1}{2}j''(\tilde{u}_k)(u_k - \bar{u})^2 \quad \text{for all } k,$$
(4.32)

where we abbreviate  $(\tilde{y}_k, \tilde{u}_k) = (\bar{y}, \bar{u}) + \gamma_k((y_k, u_k) - (\bar{y}, \bar{u}))$  with some  $\gamma_k \in [0, 1]$ . Due to  $u_k \to \bar{u}$  in  $L^r(0, T; L^2(\Omega))$ , we have  $y_k \to \bar{y}$  in  $L^{\infty}(Q)$ , see Proposition 2.4 and (2.4). Since g and j are twice continuously Fréchet-differentiable, we obtain

$$(g''(\tilde{y}_k) - g''(\bar{y})) \left(\frac{y_k - \bar{y}}{\sigma_k}\right)^2 + g''(\bar{y}) \left(\frac{y_k - \bar{y}}{\sigma_k}\right)^2 \to g''(\bar{y}) \left(S'(\bar{u};h), S'(\bar{u};h)\right) \quad \text{as } k \to \infty,$$

$$(4.33)$$

in view of (4.31), and

$$\liminf_{k \to \infty} \left( j''(\tilde{u}_k) - j''(\bar{u}) \right) \left( \frac{u_k - \bar{u}}{\sigma_k} \right)^2 + j''(\bar{u}) \left( \frac{u_k - \bar{u}}{\sigma_k} \right)^2 \ge j''(\bar{u})(h, h) \quad \text{as } k \to \infty,$$
(4.34)

in light of (4.29) and (4.24). Note that the latter one implies that  $j''(\bar{u})$  induces a norm. Dividing by  $\sigma_k$  in (4.32) and passing to the limit therein then yields

$$\lim_{k \to \infty} g'(\bar{y}) \left(\frac{y_k - \bar{y}}{\sigma_k}\right) + j'(\bar{u}) \left(\frac{u_k - \bar{u}}{\sigma_k}\right)$$

$$\leq \lim_{k \to \infty} \frac{\sigma_k}{k} - \lim_{k \to \infty} \frac{\sigma_k}{2} g''(\tilde{y}_k) \left(\frac{y_k - \bar{y}}{\sigma_k}\right)^2 + \limsup_{k \to \infty} -\frac{\sigma_k}{2} j''(\tilde{u}_k) \left(\frac{u_k - \bar{u}}{\sigma_k}\right)^2 = 0,$$
(4.35)

in view of (4.33) and (4.34). On the other hand, (4.31) and (4.29) lead to

$$\lim_{k \to \infty} g'(\bar{y}) \left(\frac{y_k - \bar{y}}{\sigma_k}\right) + j'(\bar{u}) \left(\frac{u_k - \bar{u}}{\sigma_k}\right) = g'(\bar{y}) S'(\bar{u}; h) + j'(\bar{u}) h \ge 0, \tag{4.36}$$

where the last inequality is due to (4.1) and the proofs of [23, Thm. 5.7 and Thm. 6.10]. Note that the density of the range of B in  $W^{-1,q}(\Omega)$  is not needed here. Thus, by (4.35) and (4.36), we have

$$g'(\bar{y})S'(\bar{u};h) + j'(\bar{u})h = 0.$$
(4.37)

Now, we discuss the difference of the values of the objective J. From Lemma 4.2 combined with (4.28) we know that

$$J(y_k, u_k) - J(\bar{y}, \bar{u}) \ge \int_Q p(t, x) \left( f(\bar{y}) - f(y_k) + f'(\bar{y}; y_k - \bar{y}) \right)(t, x) \, d(t, x) + \frac{1}{2} g''(\tilde{y}_k) (y_k - \bar{y})^2 + \frac{1}{2} j''(\tilde{u}_k) (u_k - \bar{u})^2 \quad \text{for all } k.$$

$$(4.38)$$

We begin by estimating the first term on the right-hand side in (4.38). To this end, we want to employ Lemmas 4.4, 4.9 and 4.12. Note that  $\inf_{z \in \mathcal{N}} \rho_z \ge \inf_{z \in \mathcal{N}} (\rho_z - \delta_z) + \inf_{z \in \mathcal{N}} \delta_z > 0$ , by Assumptions 4.6 and 4.10, so that all these three lemmas are applicable. Let now  $\varepsilon_1 > 0$  and  $\varepsilon_2 > 0$  be given by Lemmas 4.4 and 4.9, respectively. From (2.6) and (2.4) we know that there exists a constant  $C = C(\bar{u}) > 0$  such that

$$\|S(v) - \bar{y}\|_{L^{\infty}(Q)} \le C \|v - \bar{u}\|_{L^{r}(0,T;L^{2}(\Omega))}$$
(4.39)

for all  $v \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u},1)}$ . We set  $\epsilon := \min\{\varepsilon_1/2C, \varepsilon_2/2C, 1\} > 0$  and choose  $\bar{k}$  large enough such that  $u_k \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u},\epsilon)}$  for all  $k \ge \bar{k}$ . Then, by (4.39), we have  $\|S(u_k) - \bar{y}\|_{L^{\infty}(Q)} < \min\{\varepsilon_1,\varepsilon_2\}$  for all  $k \ge \bar{k}$ . We are now in the position to apply Lemmas 4.4, 4.9 and 4.12, by means of which (4.38) can be continued as

$$J(y_k, u_k) - J(\bar{y}, \bar{u}) \ge -\frac{1}{2} \int_{Q_s} p(t, x) f''(\bar{y}(t, x)) S'(\bar{u}; u_k - \bar{u})(t, x)^2 d(t, x) + o(\sigma_k^2) + \frac{1}{2} g''(\tilde{y}_k) (y_k - \bar{y})^2 + \frac{1}{2} j''(\tilde{u}_k) (u_k - \bar{u})^2 \quad \text{for all } k \ge \bar{k}.$$

$$(4.40)$$

Note that (4.40) is a result of  $Q = Q_{\mathfrak{n}} \cup Q_{\mathfrak{s},\delta} \cup Q_{\mathfrak{s}}$ , see definitions (4.3), (4.7) and (4.10). In view of (4.28) and (4.40), it further holds

$$1/k > -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'\left(\bar{u}; \frac{u_k - \bar{u}}{\sigma_k}\right) (t,x)^2 d(t,x) + \frac{o(\sigma_k^2)}{\sigma_k^2} + \frac{1}{2} g''(\tilde{y}_k) \left(\frac{y_k - \bar{y}}{\sigma_k}\right)^2 + \frac{1}{2} j''(\tilde{u}_k) \left(\frac{u_k - \bar{u}}{\sigma_k}\right)^2 \quad \text{for all } k \ge \bar{k},$$
(4.41)

where we employed the positive homogeneity of the directional derivative. We build  $\liminf_{k\to\infty}$  in (4.41), which by (4.30), (4.33) and (4.34), gives in turn

$$0 \ge -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};h)(t,x)^2 d(t,x) + \frac{1}{2} g''(\bar{y}) \left( S'(\bar{u};h), S'(\bar{u};h) \right) + \frac{1}{2} j''(\bar{u})(h,h).$$

$$(4.42)$$

By Assumption 4.13.2, which can be applied in view of (4.37), we deduce from (4.42) that h = 0. This leads to  $S'(\bar{u}; h) = 0$ , due to the definition of the directional derivative. As a result thereof, (4.30) and (4.33) now read

$$S'(\bar{u};h_k) \to 0 \quad \text{in } C([0,T];L^{\infty}(\Omega)) \quad \text{as } k \to \infty,$$
  
$$(g''(\bar{y}_k) - g''(\bar{y})) \Big(\frac{y_k - \bar{y}}{\sigma_k}\Big)^2 + g''(\bar{y}) \Big(\frac{y_k - \bar{y}}{\sigma_k}\Big)^2 \to 0 \quad \text{as } k \to \infty.$$
(4.43)

Thanks to  $\|h_k\|_{L^r(0,T;L^2(\Omega))} = 1$ , the coercivity condition (4.24), and (4.41), we obtain

$$\begin{split} 1/k &> -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'\Big(\bar{u}; \frac{u_k - \bar{u}}{\sigma_k}\Big)(t,x)^2 \, d(t,x) + \frac{o(\sigma_k^2)}{\sigma_k^2} \\ &+ \frac{1}{2} (g''(\tilde{y}_k) - g''(\bar{y})) \Big(\frac{y_k - \bar{y}}{\sigma_k}\Big)^2 + \frac{1}{2} g''(\bar{y}) \Big(\frac{y_k - \bar{y}}{\sigma_k}\Big)^2 \\ &+ \frac{1}{2} (j''(\tilde{u}_k) - j''(\bar{u})) \Big(\frac{u_k - \bar{u}}{\sigma_k}\Big)^2 + \frac{\nu}{2} \underbrace{\left\|\frac{u_k - \bar{u}}{\sigma_k}\right\|_{L^r(0,T;L^2(\Omega))}^2}_{=1} \quad \text{for all } k \ge \bar{k}. \end{split}$$

We pass again to the limit  $k \to \infty$  on both sides, which in view of (4.43), and the continuity of j'', results in  $0 \ge \nu/2 > 0$ . This finally gives the contradiction and completes the proof.  $\Box$ 

Next, we establish second-order sufficient optimality conditions for (P) which allow for an arbitrary (twice continuously Fréchet-differentiable) objective. Besides Assumptions 4.6 and 4.10, the SSC consist of a coercivity condition for the Hessian of a functional which corresponds to the classical Lagrange-functional, see Remark 4.19 below.

ASSUMPTION 4.16. In addition to Assumptions 4.6 and 4.10, we assume that there exists  $\kappa > 0$  such that

$$J''(\bar{y},\bar{u})(\eta,h)^2 - \int_{Q_s} p(t,x)f''(\bar{y}(t,x))\eta(t,x)^2 \, d(t,x) \ge \kappa \|h\|_{L^r(0,T;L^2(\Omega))}^2$$
(4.44)

for all  $h \in L^r(0,T; L^2(\Omega))$  and  $\eta = S'(\bar{u}; h)$ . Here  $Q_{\mathfrak{s}}$  denotes again the set associated with  $\{\delta_z\}_{z \in \mathcal{N}}$  given by (4.10).

REMARK 4.17. According to Assumption 4.16, the price for allowing an arbitrary objective J is the more restrictive condition (4.44). Unlike in Assumption 4.13.2, we now deal with a coercivity condition which has to be satisfied for the set of all directions, instead of the cone of critical directions. Note that, in contrast to finite dimensions, (4.44) is not necessarily equivalent to

$$J''(\bar{y},\bar{u})(\eta,h)^2 - \int_{Q_s} p(t,x) f''(\bar{y}(t,x))\eta(t,x)^2 d(t,x) > 0$$

for all  $h \in L^r(0,T; L^2(\Omega)) \setminus \{0\}$  and  $\eta = S'(\bar{u}; h)$ , see [10, Section 3.2]. However, when considering the objective  $\hat{J}(y, u) = \frac{1}{2} ||y - y_d||^2_{L^2(Q)} + \frac{\nu}{2} ||u - u_d||^2_{H^1(0,T; L^2(\Omega))}$ and the control space as in Remark 4.14.(i), this equivalence holds. To see this, one employs the same arguments as in the finite dimensional case combined with the compact embedding (2.4). We refer here also to the proof of [10, Thm 4.11]. If  $J = \hat{J}$ , (4.44) is satisfied when

$$p(t,x)f''(\bar{y}(t,x)) \le 1$$
 a.e. in  $Q_{\mathfrak{s}}$ , (4.45)

thanks to the embedding  $H^1(0,T; L^2(\Omega)) \hookrightarrow L^{\infty}(0,T; L^2(\Omega))$ . In the next section, we provide conditions on the given data which ensure (4.45).

THEOREM 4.18. Let  $(\bar{u}, \bar{y}, \lambda, p)$  satisfy the first-order optimality system (4.1) given by Theorem 4.1. If Assumptions 4.6, 4.10 and 4.16 are fulfilled, then there exist  $\alpha > 0$ and R > 0 such that

$$J(\bar{y},\bar{u}) + \alpha \|u - \bar{u}\|_{L^{r}(0,T;L^{2}(\Omega))}^{2} \leq J(S(u),u) \quad \forall u \in B_{L^{r}(0,T;L^{2}(\Omega))}(\bar{u},R).$$
(4.46)

In particular,  $\bar{u}$  is locally optimal for (P).

*Proof.* In the proof of Theorem 4.15 we already checked that, under Assumptions 4.6 and 4.10, Lemma 4.4 is applicable (in addition to Lemmas 4.9 and 4.12). We define again  $\epsilon := \min\{\varepsilon_1/2C, \varepsilon_2/2C, 1\} > 0$  and fix  $u \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u}, \epsilon)}, u \neq \bar{u}$ , where  $\varepsilon_1 > 0$  and  $\varepsilon_2 > 0$  are given by Lemmas 4.4 and 4.9, respectively, and  $C = C(\bar{u}) > 0$ is a constant so that (4.39) holds. Then, one has  $\|y - \bar{y}\|_{L^{\infty}(Q)} < \min\{\varepsilon_1, \varepsilon_2\}$ , see the proof of Theorem 4.15, where y := S(u) from now on. From Lemma 4.2 combined with Lemmas 4.4, 4.9, and 4.12 we have

$$J(y,u) - J(\bar{y},\bar{u}) \ge -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};u-\bar{u})(t,x)^2 d(t,x) + o(||u-\bar{u}||^2_{L^r(0,T;L^2(\Omega))}) + \frac{1}{2} J''(y_{\gamma_u},u_{\gamma_u})(y-\bar{y},u-\bar{u})^2,$$

$$(4.47)$$

where  $(y_{\gamma_u}, u_{\gamma_u}) = (\bar{y}, \bar{u}) + \gamma_u((y, u) - (\bar{y}, \bar{u}))$  with some  $\gamma_u \in [0, 1]$ . Further, by means of some algebraic manipulations, see the proof of [3, Thm. 4.6], we can write

$$J''(y_{\gamma_{u}}, u_{\gamma_{u}})(y - \bar{y}, u - \bar{u})^{2} - J''(\bar{y}, \bar{u})(S'(\bar{u}; u - \bar{u}), u - \bar{u})^{2}$$

$$= (J''(y_{\gamma_{u}}, u_{\gamma_{u}}) - J''(\bar{y}, \bar{u}))(y - \bar{y}, u - \bar{u})^{2} + \partial_{y}^{2}J(\bar{y}, \bar{u})(y - \bar{y} - S'(\bar{u}; u - \bar{u}))^{2}$$

$$+ 2\partial_{y}^{2}J(\bar{y}, \bar{u})(y - \bar{y} - S'(\bar{u}; u - \bar{u}), S'(\bar{u}; u - \bar{u}))$$

$$+ 2\partial_{y}\partial_{u}J(\bar{y}, \bar{u})(y - \bar{y} - S'(\bar{u}; u - \bar{u}), u - \bar{u})$$

$$= o(||u - \bar{u}||_{L^{r}(0,T;L^{2}(\Omega))}^{2}).$$
(4.48)

The last equality is due to Assumption 2.1.7 combined with Proposition 2.4, the Bouligand-differentiability of S, cf. Theorem 3.2, and Theorem 2.5 (which tells us that  $\|S'(\bar{u}; u - \bar{u})\|_{L^r(0,T;W_D^{1,q}(\Omega))} \leq c \|u - \bar{u}\|_{L^r(0,T;L^2(\Omega))}$ ). Inserting (4.48) in (4.47) and employing Assumption 4.16 results in

$$\begin{split} J(S(u),u) - J(\bar{y},\bar{u}) &\geq -\frac{1}{2} \int_{Q_s} p(t,x) f''(\bar{y}(t,x)) S'(\bar{u};u-\bar{u})(t,x)^2 \, d(t,x) \\ &+ \frac{1}{2} J''(\bar{y},\bar{u}) (S'(\bar{u};u-\bar{u}),u-\bar{u})^2 + o(\|u-\bar{u}\|_{L^r(0,T;L^2(\Omega))}^2) \\ &\geq \Big(\frac{\kappa}{2} - \frac{|o(\|u-\bar{u}\|^2)|}{\|u-\bar{u}\|_{L^r(0,T;L^2(\Omega))}^2} \Big) \|u-\bar{u}\|_{L^r(0,T;L^2(\Omega))}^2 \end{split}$$

for all  $u \in \overline{B_{L^r(0,T;L^2(\Omega))}(\bar{u},\epsilon)}$  with  $u \neq \bar{u}$ . The proof is now complete.  $\Box$ 

REMARK 4.19. If the nonlinearity  $f : \mathbb{R} \to \mathbb{R}$  is twice continuously differentiable, then both sets of SSC derived in this section coincide with the classical second-order sufficient optimality conditions, as explained in what follows. In finite dimensions, it is well known that the SSC consist of strong stationarity (necessary conditions for local optimality) and the coercivity/positive-definiteness of the Hessian (w.r.t. the primal variables) of the Lagrangian on the cone of critical directions, see e.g. [29]. For the treatment of SSC in infinite dimensions, we refer to [10]. Let us see what happens when  $f : \mathbb{R} \to \mathbb{R}$  is twice continuously differentiable in our case too. Then,  $\mathcal{N} = \emptyset$  and  $Q_s = Q$ , and thus, Assumptions 4.6 and 4.10 are automatically fulfilled. In particular, f does not have to be local convex nor local concave and the sign of p is no longer an issue. Moreover, (4.25) reads

$$g''(\bar{y})(\eta,\eta) + j''(\bar{u})(h,h) - \int_{Q} p(t,x)f''(\bar{y}(t,x))\eta(t,x)^{2} d(t,x) > 0$$
(4.49)

for all  $h \in L^r(0,T;L^2(\Omega)) \setminus \{0\}$  and  $\eta = S'(\bar{u};h)$  with  $g'(\bar{y})\eta + j'(\bar{u})h = 0$ , i.e.,  $\partial^2_{(y,u)} \mathfrak{L}(\bar{y},\bar{u},p)(\eta,h)^2 > 0$ , where the Lagrangian  $\mathfrak{L} : \mathbb{W}^r_0(W^{1,q}_D(\Omega), W^{-1,q}(\Omega)) \times L^r(0,T;L^2(\Omega)) \times \mathbb{W}^{r'}_T(W^{1,q'}_D(\Omega), W^{-1,q'}(\Omega)) \to \mathbb{R}$  is given by

$$\mathfrak{L}(y, u, p) = J(y, u) - \langle p, \dot{y} + Ay + f(y) - Bu \rangle_{L^{r}(0, T; W^{-1, q}(\Omega))}.$$
19

In conclusion, for Theorem 4.15 to hold, one only needs strong stationarity, Assumption 4.13.1 and (4.49). Let us point out that the structural assumption on J is to be expected, since this is essential in infinite dimensions in order to obtain a contradiction at the end of the proof of Theorem 4.15, see also [10] (smooth case with control constraints) and [3, 21] (non-smooth case). If f is twice continuously differentiable, the assertion in Theorem 4.18 corresponds entirely to the finite case: if  $(\bar{u}, \bar{y}, \lambda, p)$  is strong stationary and if the coercivity condition

$$\partial_{(y,u)}^2 \mathfrak{L}(\bar{y}, \bar{u}, p)(\eta, h)^2 \ge \kappa \|h\|_{L^r(0,T; L^2(\Omega))}^2$$

holds for all  $h \in L^r(0,T; L^2(\Omega))$  and  $\eta = S'(\bar{u}; h)$ , then  $\bar{u}$  is a strict local optimum which fulfills (4.46). This coincides with the assertion in [6, Thm. 8.3.3], which deals with the smooth case in infinite dimensions.

5. Second-order sufficient conditions for a concrete setting. In this section, we derive conditions on the given data under which Assumption 4.6.2-3 and (4.45) are guaranteed. To this end, we consider the optimal control problem

$$\begin{array}{ccc} \min_{u \in H^{1}(0,T;L^{2}(\Omega))} & \frac{1}{2} \|y - y_{d}\|_{L^{2}(Q)}^{2} + \frac{1}{2} \|u\|_{H^{1}(0,T;L^{2}(\Omega))}^{2} \\ \text{s.t.} & \dot{y}(t) - \Delta y(t) + f(y(t)) = u(t) \quad \text{a.e. in } (0,T), \\ & y(0) = 0. \end{array} \right\}$$

$$(P_{ex})$$

The fact that we choose  $H^1(0, T; L^2(\Omega))$  as the control space does not affect the previous results, see Remark 4.14.(i). In all what follows,  $B := L^2(\Omega) \hookrightarrow W^{-1,q}(\Omega)$ is the embedding operator. This is well-defined, provided that  $q \leq 2n/(n-2)$ . In  $(P_{ex}), -\Delta : W_D^{1,q}(\Omega) \to W^{-1,q}(\Omega)$  denotes the Laplace operator in the distributional sense, i.e.,  $-\Delta = -\operatorname{div} \nabla$ . We assume that  $\Omega$  is smooth enough and  $\partial\Omega = \Gamma_D$ , so that Assumptions 2.1.2-3 are satisfied by  $-\Delta$ , cf. Remark 2.7. Note that  $\Omega$  is regular in the sense of Gröger, in view of [16, Thm. 5.2 and 5.4]. The desired state  $y_d$  belongs to  $L^r(0,T; L^2(\Omega))$ , while r is supposed to fulfill (2.1) where  $q \in (n, 2n/(n-2)]$  is fixed. In the following, we abbreviate

$$p^+ := \max\{p, 0\}$$
 and  $p^- := \min\{p, 0\}$ .

By  $C_{\Omega}$  we denote the Poincaré constant associated with the domain  $\Omega$ .

Throughout this section,  $(\bar{u}, \bar{y}, \lambda, p)$  is a fixed point that satisfies the first-order optimality system given by Theorem 4.1. In the setting considered here, (4.1) reads

$$\dot{\bar{y}} - \bigtriangleup \bar{y} + f(\bar{y}) = \bar{u}, \quad \bar{y}(0) = 0, \tag{5.1a}$$

$$-\dot{p} - \triangle p + \lambda = \bar{y} - y_d, \quad p(T) = 0, \tag{5.1b}$$

$$\lambda(t,x) \in [f'_{+}(\bar{y}(t,x)) \, p(t,x), f'_{-}(\bar{y}(t,x)) \, p(t,x)] \text{ a.e. in } Q, \tag{5.1c}$$

$$p + \bar{u} = 0. \tag{5.1d}$$

We suppose that the nonlinearity satisfies f(0) = 0 and that f is convex around any  $z \in \mathcal{N}$  with radius  $\rho_z > 0$  (in addition to Assumptions 2.1.4-5). Then,  $f(y)y \ge 0$  and  $f'(y;h)h \ge 0$  for all  $y, h \in \mathbb{R}$ . Thus, by (5.1c), we find

$$f(\bar{y}(t,x))\bar{y}(t,x) \ge 0, \quad \lambda(t,x)p(t,x) \ge 0, \ \lambda(t,x)p^+(t,x) \ge 0$$
 a.e. in  $Q.$  (5.2)  
20



FIG. 5.1. An ignition-type nonlinearity with ignition temperatures  $\theta_1, \theta_2$ , and  $\theta_3$ 

For the above-described problem we show that if (5.9) below is satisfied, then  $p(t, x) \leq 0$  a.e. in Q, i.e., Assumption 4.6.2 holds true. Let us define  $\delta_z := \rho_z/2$  for all  $z \in \mathcal{N}$  in Assumption 4.6. If, in addition to (5.9) below,  $\mathcal{N}$  is finite, f is twice continuously differentiable on  $\{v \in \mathbb{R} | |v-z| \geq \rho_z/4 \ \forall z \in \mathcal{N}\}$ , and (5.15) below holds, we can also prove (4.45). Thus, we provide a setting for which all assumptions in Theorems 4.15 and 4.18 are satisfied, see also Remarks 4.11, 4.14, and 4.17. In this context, (5.1) is not only necessary but also sufficient for optimality.

Before we begin with the proof, we mention that a nonlinearity which satisfies all the above conditions (i.e., Lipschitz continuous on bounded sets, directionally differentiable with  $\mathcal{N}$  finite, convex around any  $z \in \mathcal{N}$ , monotone increasing with f(0) = 0 and twice continuously differentiable on a subset of  $\mathcal{S}$ ) is depicted in Figure 5.1. This could arise in combustion processes where different ignition temperatures are given, see Remark 2.9.

(I) We first deal with Assumption 4.6.2. We start by showing that  $||p||_{L^2(Q)} \leq K$ , where K > 0 is some constant which depends only on the given data (step (i) below). This will enable us to derive conditions on the data such that  $\bar{y} \leq y_d$  a.e. in Q holds (step (ii) below). By means of this inequality, we can then conclude that  $p \leq 0$  a.e. in Q (step (iii) below).

(i) As a consequence of (5.1d), we have  $p \in H^1(0,T;L^2(\Omega))$ , which gives in turn  $\lambda \in L^{\infty}(0,T;L^2(\Omega))$ , by (5.1c) and (3.2). Thus,  $t \mapsto (\bar{y}-y_d-\lambda)(T-t) \in L^r(0,T;L^2(\Omega)) \hookrightarrow L^r(0,T;W^{-1,q}(\Omega))$ , since  $q \leq 2n/(n-2)$ . Since  $-\Delta$  satisfies maximal parabolic  $L^r(0,T;W^{-1,q}(\Omega))$ -regularity, we now deduce from (5.1b) that  $p \in W^r_0(W^{1,q}_D(\Omega),W^{-1,q}(\Omega))$ , where we used the transformation  $t \mapsto T-t$ . Therefore, we can test the adjoint equation (5.1b) with  $p(T-\cdot)$ , which leads to

$$\frac{1}{2} \|p(T-t)\|_{L^{2}(\Omega)}^{2} - \frac{1}{2} \underbrace{\|p(T)\|_{L^{2}(\Omega)}^{2}}_{=0} + \frac{1}{C_{\Omega}^{2}} \int_{0}^{t} \|p(T-s)\|_{L^{2}(\Omega)}^{2} ds$$
$$\leq \int_{0}^{t} \int_{\Omega} (\bar{y} - y_{d})(T-s, x) p(T-s, x) d(s, x) \quad \forall t \in [0, T]$$
(5.3)

in view of the formula of integration by parts, Poincaré-Friedrichs's inequality and (5.2). We proceed in the same way regarding the state equation. We test (5.1a) with y and employ the first inequality in (5.2), as well as Poincaré-Friedrichs's inequality.

The resulting estimate combined with (5.1d) is then used on the right-hand side of (5.3) at t := T, which leads to

$$\frac{1}{C_{\Omega}^{2}} \int_{0}^{T} \|p(s)\|_{L^{2}(\Omega)}^{2} ds \leq \underbrace{\|\bar{y}(0)\|_{L^{2}(\Omega)}^{2}}_{=0} - \frac{1}{C_{\Omega}^{2}} \int_{0}^{T} \|\bar{y}(s)\|_{L^{2}(\Omega)}^{2} ds - \int_{Q} y_{d}(s,x)p(s,x) d(s,x)$$
(5.4)

From (5.4) we deduce

$$2\|p\|_{L^{2}(Q)}\|\bar{y}\|_{L^{2}(Q)} \leq \|p\|_{L^{2}(Q)}^{2} + \|\bar{y}\|_{L^{2}(Q)}^{2} \leq C_{\Omega}^{2}\|y_{d}\|_{L^{2}(Q)}\|p\|_{L^{2}(Q)}.$$
(5.5)

Dividing by  $||p||_{L^2(Q)}$  results in

$$\|p\|_{L^{2}(Q)} \leq C_{\Omega}^{2} \|y_{d}\|_{L^{2}(Q)}, \quad \|\bar{y}\|_{L^{2}(Q)} \leq \frac{C_{\Omega}^{2} \|y_{d}\|_{L^{2}(Q)}}{2}.$$
(5.6)

(ii) In order to prove that  $p \leq 0$ , we derive conditions which guarantee that  $\bar{y} \leq y_d$  a.e. in Q. To this end, we insert (5.5) on the right-hand side in (5.3) and obtain

$$\|p(T-t)\|_{L^{2}(\Omega)}^{2} \leq (C_{\Omega}^{2}+2)\|y_{d}\|_{L^{2}(Q)}\|p\|_{L^{2}(Q)} \leq (C_{\Omega}^{2}+2)C_{\Omega}^{2}\|y_{d}\|_{L^{2}(Q)}^{2} \quad \forall t \in [0,T],$$
(5.7)

where for the last inequality we used (5.6). Via a comparison principle, cf. [27, Lem. A.1, Prop. 3.3 and 3.4], where one relies on the monotonicity of f, one can show that

$$|\bar{y}(t,x)| \le \tilde{y}(t,x) \le \underbrace{C_e \| (\partial_t - \Delta)^{-1} \|_{\mathcal{L}(L^r(0,T;W^{-1,q}(\Omega)), \mathbb{W}_0^r(W_D^{1,q}(\Omega), W^{-1,q}(\Omega)))}}_{=:K_1} \| p \|_{L^r(0,T;L^2(\Omega))} \quad \text{a.e. in } Q,$$

(5.8) where  $\tilde{y} \in \mathbb{W}_{D}^{r}(W_{D}^{1,q}(\Omega), W^{-1,q}(\Omega))$  is the unique solution of  $\dot{\tilde{y}} - \Delta \tilde{y} = |\bar{u}|$  and  $C_{e} > 0$ is the product of the embedding constants of  $\mathbb{W}_{0}^{r}(W_{D}^{1,q}(\Omega), W^{-1,q}(\Omega)) \hookrightarrow L^{\infty}(Q)$  and  $L^{2}(\Omega) \hookrightarrow W^{-1,q}(\Omega)$ , cf. (2.4). Note that for (5.8) we employed again (5.1d). Thus, in view of (5.7) and (5.8), we deduce that if

$$K_1 T^{1/r} \sqrt{C_{\Omega}^2 + 2C_{\Omega} \|y_d\|_{L^2(Q)}} \le y_d \text{ a.e. in } Q,$$
 (5.9)

then  $\bar{y} \leq y_d$  a.e. in Q.

(iii) Now, to see that (5.9) implies the desired result, we test the equation (5.1b) with  $p^+ \in L^r(0,T; W_D^{1,q}(\Omega))$ , see [19, Thm. A.1]. We arrive at

$$\int_{0}^{t} \langle -\dot{p}(T-s), p^{+}(T-s) \rangle_{W_{D}^{1,q}(\Omega)} ds + \int_{0}^{t} \int_{\Omega} \underbrace{\nabla p(T-s)(x) \nabla p^{+}(T-s)(x)}_{\geq 0} d(s,x) + \int_{0}^{t} \int_{\Omega} \underbrace{\lambda(T-s,x) p^{+}(T-s,x)}_{\geq 0, \text{ see } (5.2)} d(s,x) = \int_{0}^{t} \int_{\Omega} \underbrace{(\bar{y}-y_d)(T-s,x)}_{\leq 0} p^{+}(T-s,x) d(s,x) \quad \forall t \in [0,T].$$
(5.10)

Thanks to [32, Lemma 3.2] combined with (5.10) and  $p \in H^1(0,T; L^2(\Omega))$ , we have

$$\int_0^t \langle -\dot{p}(T-s), p^+(T-s) \rangle_{W_D^{1,q}(\Omega)} \, ds = 1/2 \| p^+(T-t) \|_{L^2(\Omega)}^2 - 1/2 \underbrace{\| p^+(T) \|_{L^2(\Omega)}^2}_{=0} \le 0 \quad \text{for all } t \in [0,T],$$

which gives in turn  $p \leq 0$  a.e. in Q.

(II) In order to derive conditions which guarantee (4.45), we first show that there exists a constant c > 0, dependent only on the given data, so that  $-p \leq c$  a.e. in Q. To this end, we apply a comparison principle again and assume in the following that (5.9) holds. Consider the equation

$$-\dot{\tilde{p}} - \Delta \tilde{p} = \bar{y} - y_d, \quad \tilde{p}(T) = 0.$$
(5.11)

Since  $-\triangle$  satisfies maximal parabolic  $L^r(0,T; W^{-1,q}(\Omega))$ -regularity, there exists a unique  $\tilde{p} \in W^r_T(W^{1,q}_D(\Omega), W^{-1,q}(\Omega)) \hookrightarrow L^{\infty}(Q)$  which solves (5.11). To see this, one uses the transformation  $t \mapsto T - t$ . Thus,

$$-\tilde{p}(t,x) \le \|\tilde{p}\|_{L^{\infty}(Q)} \le K_1 \|\bar{y} - y_d\|_{L^r(0,T;L^2(\Omega))} \le c \quad \text{a.e. in } Q,$$
(5.12)

where  $K_1$  denotes the constant from (5.8). Note that a value for c (dependent only on the given data) can be obtained from (5.7) and (5.8). In view of (5.2) and  $p \leq 0$ ,  $\lambda \leq 0$  follows. We "compare" (5.1b) with (5.11) and we see that

$$-p(t,x) \le -\tilde{p}(t,x) \le c \quad \text{a.e. in } Q, \tag{5.13}$$

as a result of (5.12). Here we relied on [32, Lemma 3.3]. Before we proceed with the proof, let us recall that we defined  $\delta_z := \rho_z/2$  for all  $z \in \mathcal{N}$ , which means that

$$Q_{\mathfrak{s}} = \{(t, x) \in Q \mid |\bar{y}(t, x) - z| \ge \rho_z/2 \ \forall z \in \mathcal{N}\}.$$

Note that (4.45) is automatically satisfied a.e. in  $\{(t,x) \in Q_{\mathfrak{s}} | f''(\bar{y}(t,x)) \ge 0\}$ , since  $p \le 0$ . On the other hand, from (5.13) one has

$$\begin{aligned}
f''(\bar{y}(t,x))p(t,x) &\leq -cf''(\bar{y}(t,x)) = c|f''(\bar{y}(t,x))| \\
&\leq c \sup_{v \in \mathcal{M}} |f''(v)| \\
&< \infty \quad \text{f.a.a.} \ (t,x) \in Q_{\mathfrak{s}} \text{ with } f''(\bar{y}(t,x)) < 0,
\end{aligned} \tag{5.14}$$

where we abbreviate  $\mathcal{M} := \{v \in \mathbb{R} \mid |v-z| \ge \rho_z/2 \,\forall z \in \mathcal{N}, |v| \le \|\bar{y}\|_{L^{\infty}(Q)}, f''(v) < 0\}$ . The last inequality in (5.14) is true, since  $\mathcal{N}$  is finite and f'' is continuous on  $\{v \in \mathbb{R} \mid |v-z| \ge \rho_z/4 \ \forall z \in \mathcal{N}\}$ , by assumption. Thus, if in addition to (5.9),

$$c \sup_{v \in \mathcal{M}} |f''(v)| \le 1 \tag{5.15}$$

holds, then (4.45) is guaranteed. Note that (4.45) is automatically fulfilled if f is convex on  $\{v \in \mathbb{R} \mid |v-z| \ge \rho_z/4 \,\forall z \in \mathcal{N}, |v| \le \|\bar{y}\|_{L^{\infty}(Q)} + 1\}$ , since in this case  $f''(\bar{y}(t,x)) \ge 0$  f.a.a.  $(t,x) \in Q_{\mathfrak{s}}$ .

In conclusion, for the setting considered here, (5.9) and (5.15) imply that every strong stationary point of (P) satisfies all assumptions in Theorems 4.15 and 4.18, and thus, the necessary optimality condition (5.1) is also sufficient for local optimality.

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